Statistical Mechanics in the Galactic Center Mass segregation and phase transitions

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Young Stars in the Galactic center



Figure: Overdensity of angular momentum distribution of the young stars in log scale at different radius slices².

• More massive young stars form a coeval warped disk, with a possible counterrotating structure.

¹Sebastiano D. von Fellenberg et al. "The Young Stars in the Galactic Center". In: ApJ 932.1, L6 (June 2022), p. L6. DOI: 10.3847/2041-8213/ac68ef. arXiv: 2205.07595 [astro-ph.GA]. ²von Fellenberg et al., "The Young Stars in the Galactic Center".

Vector Resonant Relaxation



- Averaged over orbit and precession.
- Fixed semi major axis and eccentricity.
- Efficiently change the inclinations of the stars. ³

³Figures adopted from Kinetic theory notes on self-gravitating systems, Fouvry, 2022.

The main objectives of this presentation are:

- A statistical mechanical toy model for stars in the galactic center under vector resonant relaxtion (VRR);
- Heavier stars condense into a disk: mass segregation in the inclination of the stars;
- Order-disorder phase transition;
- Negative absolute temperatures;
- Compare thermal equilibrium with dynamical simulations;

Statistical Mechanics Models

Two-component Models

• Nuclear star clusters with two components denoted by C_1 and C_2 of N_1 and N_2 number of stars, respectively, orbiting around the same central SMBH, around the same axis of symmetry.



Two-component Models

• The VRR Hamiltonian with quadrupole approximation is given by

$$H_{\text{VRR}} = -\frac{1}{2} \sum_{i,j \in \mathcal{C}_1} J_1 P_2(\boldsymbol{n}_i \cdot \boldsymbol{n}_j) - \frac{1}{2} \sum_{i,j \in \mathcal{C}_2} J_2 P_2(\boldsymbol{n}_i \cdot \boldsymbol{n}_j) - \sum_{i \in \mathcal{C}_1, j \in \mathcal{C}_2} J' P_2(\boldsymbol{n}_i \cdot \boldsymbol{n}_j),$$
(1)

- *n_i* is the unit angular momentum vector direction of the *i*-the star.
- *P*₂ is the second Legendre polynomials
- $J_1, J_2/J'$ are coupling constants.

The equilibrium distribution function

- Extremising the total entropy subject to fixed total energy and total angular momentum.
- Using mean-field approximation, the equilibrium distribution function of the inclination *i* is given by:

$$f_1(i|Q_1,Q_2) \propto e^{\frac{3}{2}\beta(J_1N_1Q_1+J'N_2Q_2)\cos^2 i+l_1\gamma\cos i}.$$
 (2)

- Lagrange multipliers: β (temperature), γ (net rotation rate).
- Quadrupole moment

$$Q_{1,2} = \langle \cos^2 i - 1/3 \rangle \tag{3}$$

The order parameter Q



• Maximum value Q = 2/3, razor-thin disk, ordered. ^a

^ahttps://galaxiesbook.org/chapters/II-01.-Flattened-Mass-Distributions.html



• Q = 0, isotropic distribution, disordered. ^a

^ahttps://sci.esa.int/web/hubble/-/globular-cluster-ngc-1466

Anisotropic Mass Segregation

The heat bath approximation

- One of the two components dominates the total energy and total angular momentum, acting as a heat bath, a reservoir of energy and angular momentum.
- The subdominant star cluster then determines its equilibrium state according to the temperatures β, γ set by the heat bath.
- Similar to the canonical ensemble



Vertical Mass Segregation

- Fix the distribution of the heat bath and vary the mass ratio of individual stars to the heat bath.
- Observe the change in the order parameter Q of the subdominant component, Wang & Kocsis 2023⁴.
- Q close to 0: isotropic sphere, Q close to 2/3: disk-like.
- Heavier stars, larger Q values, condense into a disk



⁴Hanxi Wang and Bence Kocsis. "Anisotropic mass segregation: Two-component mean-field model". In: Phys. Rev. D 108.10, 103004 (Nov. 2023), p. 103004. DOI: 10.1103/PhysRevD.108.103004. arXiv: 2302.12842 [astro-ph.GA].

Dynamical Simulation Results

- N-body simulations of the VRR Hamiltonian.
- Fix the heat bath and total star number ratio $(N_1/N_2 = 30)$.
- Start with a nearly isotropic subdominant component and let it evolve. 3100 total stars were used.
- Time-averaging the final states.



Phase Transition

Phase transition in the microcanonical ensemble

- A general two-component system (no heat bath).
- System will tend to the global maximum of entropy out of all equilibrium states $\delta S = 0$.
- Transition between water (Q near 0, disordered) and ice (Q large, ordered disk), Wang & Kocsis, in prep.



Throwing ice in water.



Water becomes ice as well.

- We evaluate the time-averaged order parameter Q as time evolves in the simulations.
- Q value increases from the sphere value to the disk value.



Initial and final states

- Initially a spherical state
- Final state follows the thermal equilibrium of the disk.







• Changes between disk and spherical states.



Negative Temperature States

Negative absolute temperature states

- Energy must be bounded.
- Entropy decreases when energy increases. $1/T \sim dS/dE$
- Boltzmann distribution inverted.
- $\bullet\,$ Observed in magnetic systems, cold atoms in optical lattices, laser systems etc. 5



⁵Figure from https://www.quantum-munich.de/119947/Negative-Absolute-Temperatures

Negative temperature states of stellar orbits.

- Bounded total energy for VRR, negative temperature states allowed. $1/T \sim dS/dE$
- Higher total angular momentum systems have more negative temperature states. (Roupas et al. 2017, Wang & Kocsis 2023)
- disordered and spherical, more energetic microstates relatively more populated.

•
$$N_1/N_2 = 15, m_1/m_2 = 0.5, a_1 = a_2, L_{\text{norm}} = 0.57.$$



Negative temperature state simulations

- Component 2 changes from a positive temperature to negative temperature
- Component 1 keeps its negative temperature.



Negative temperatures in the galactic center?



- We have built a toy model of the statistical mechanics of star orbits in the galactic center.
- The thermal equilibrium under VRR displays a mass segregation in the inclination. Higher-mass stars condense into a flattened disk under the heat bath.
- We explore the possibility of phase transition: the coexistence of 'water' and 'ice'.
- We demonstrate gravitating systems with negative temperatures.
- The dynamical simulations show good agreement with the thermal theory.