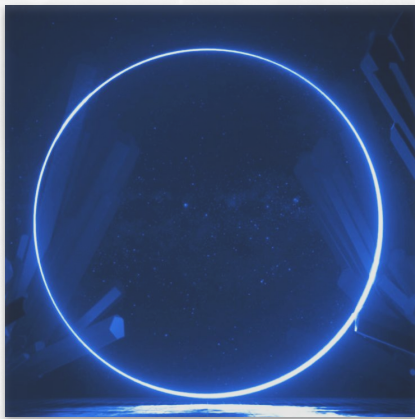




Spin estimations from lensing of hot spot around SMBHs

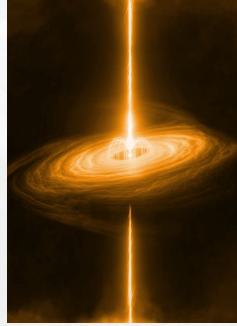
From Empiric to Deep Learning

Aristomenis Yfantis

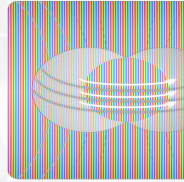
+R. Al-Belmeisi, D. Palumbo, M. Moscibrodzka

Why

➤ BH creation

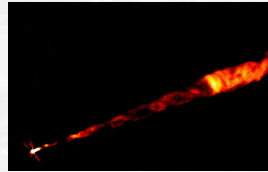


➤ Evolution

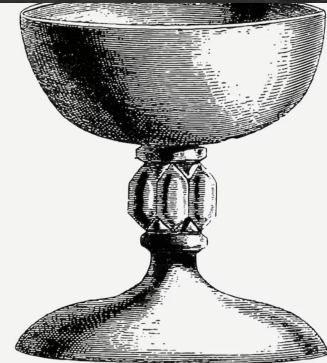
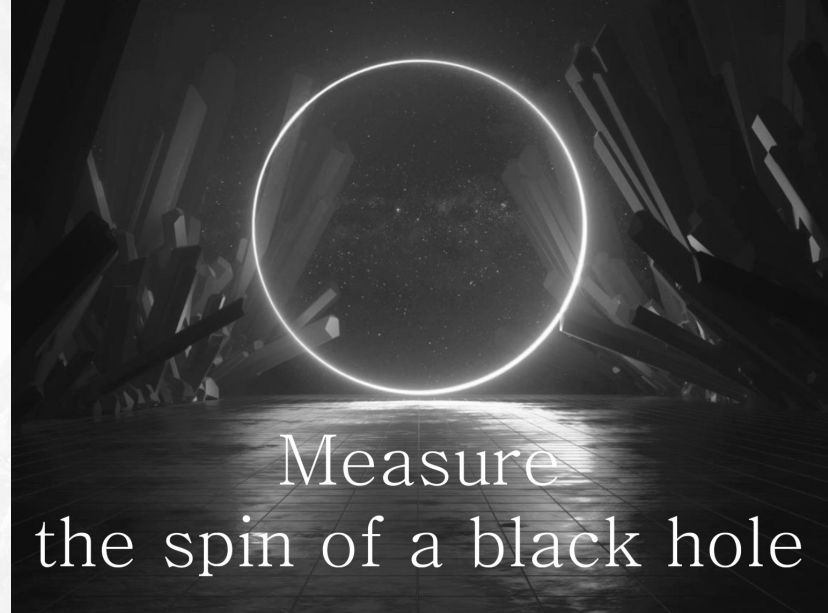
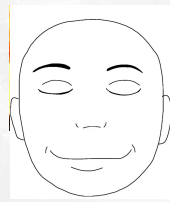


➤ Ergosphere

➤ Jets

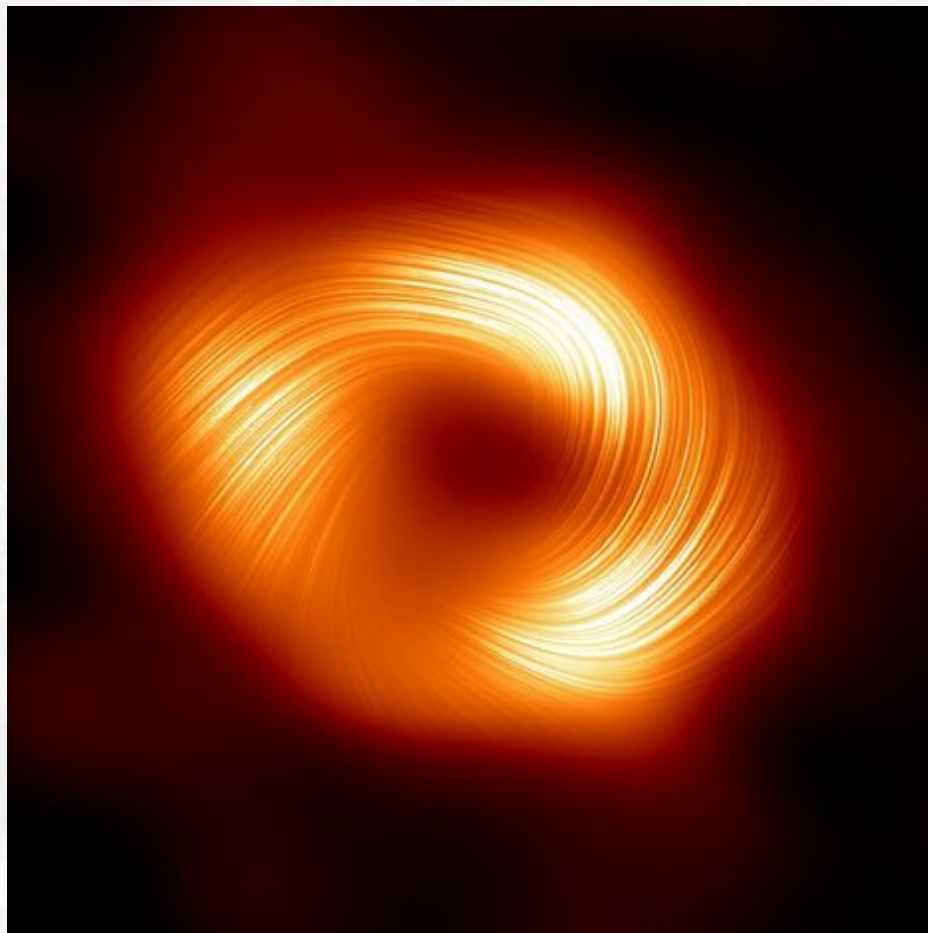


➤ Its cool

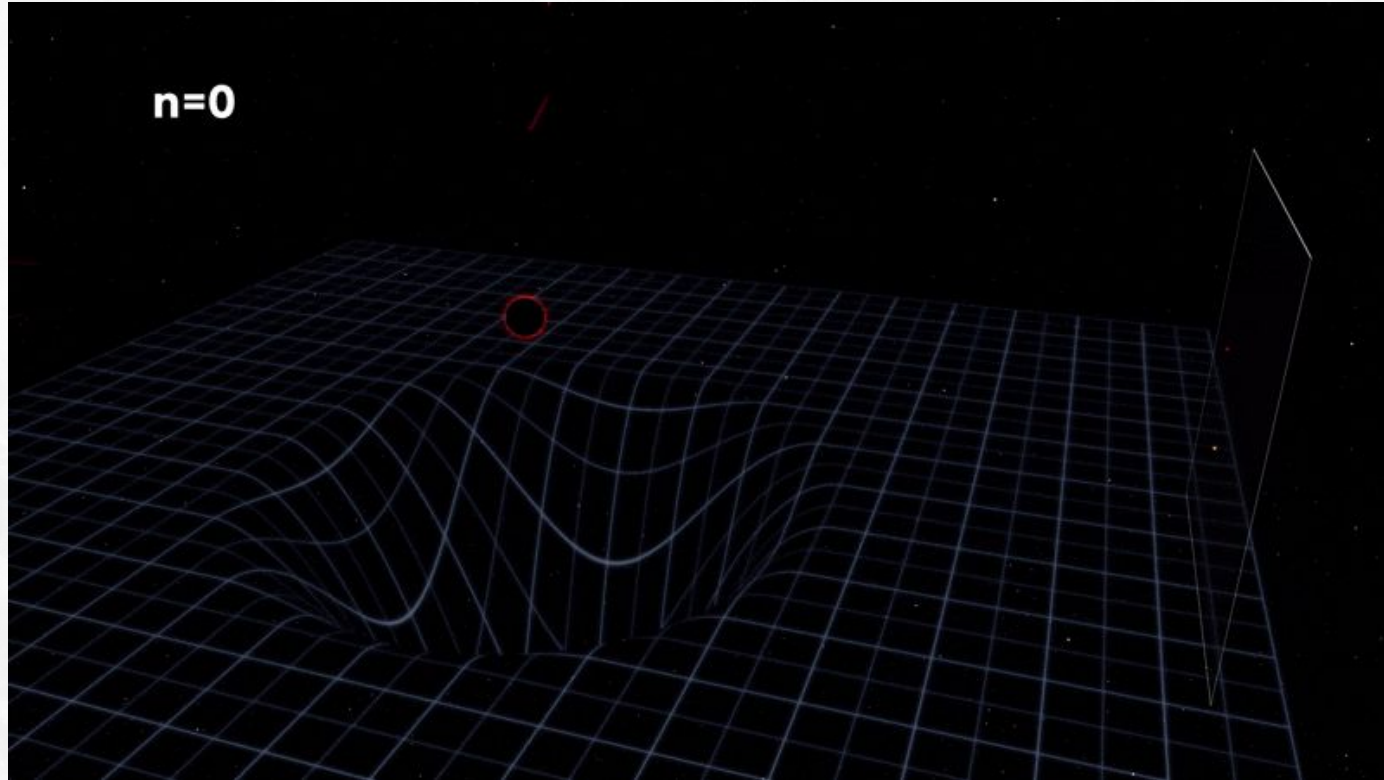


Sgr A* polarized image

- Spin information hidden somewhere in there
- But deeply mixed in the physics of accretion (dirty astrophysics)



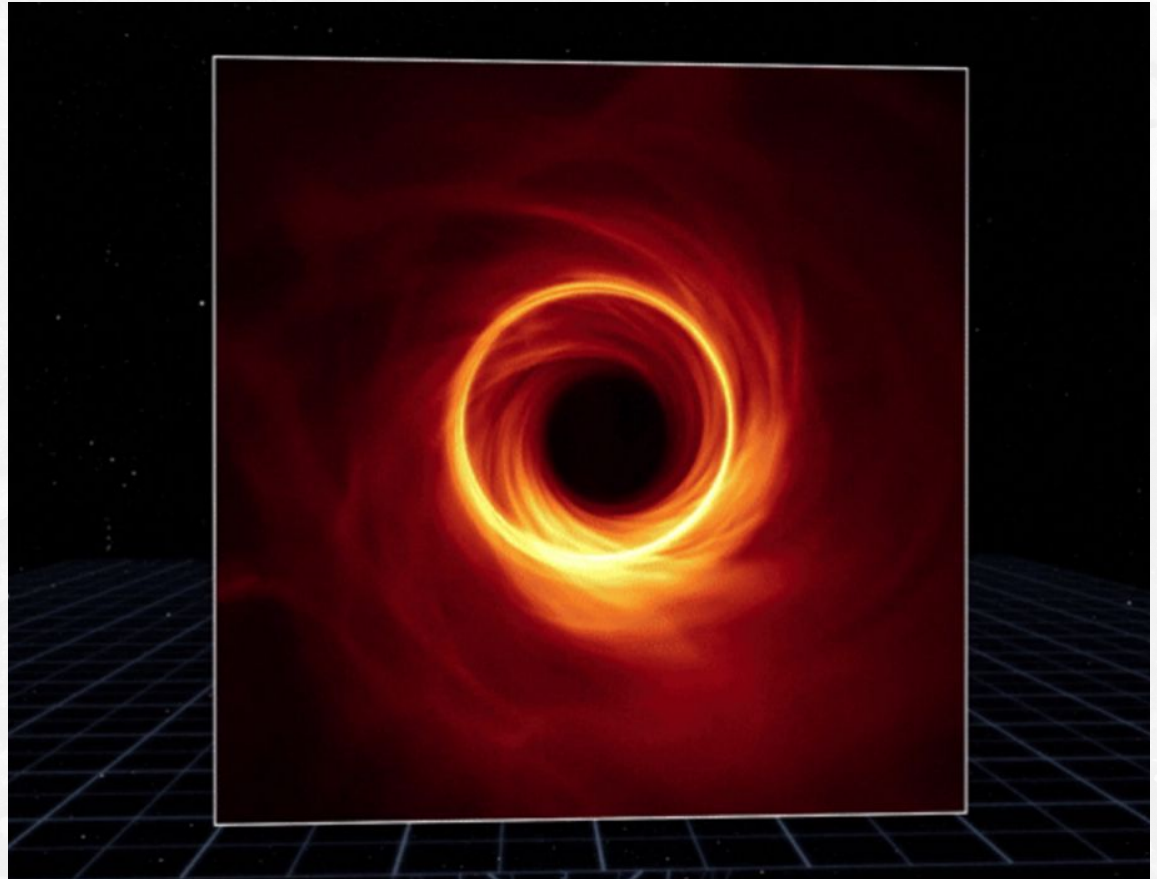
Secondary image: the photon ring



Credit: Center for Astrophysics | Harvard & Smithsonian

Secondary image: the photon ring

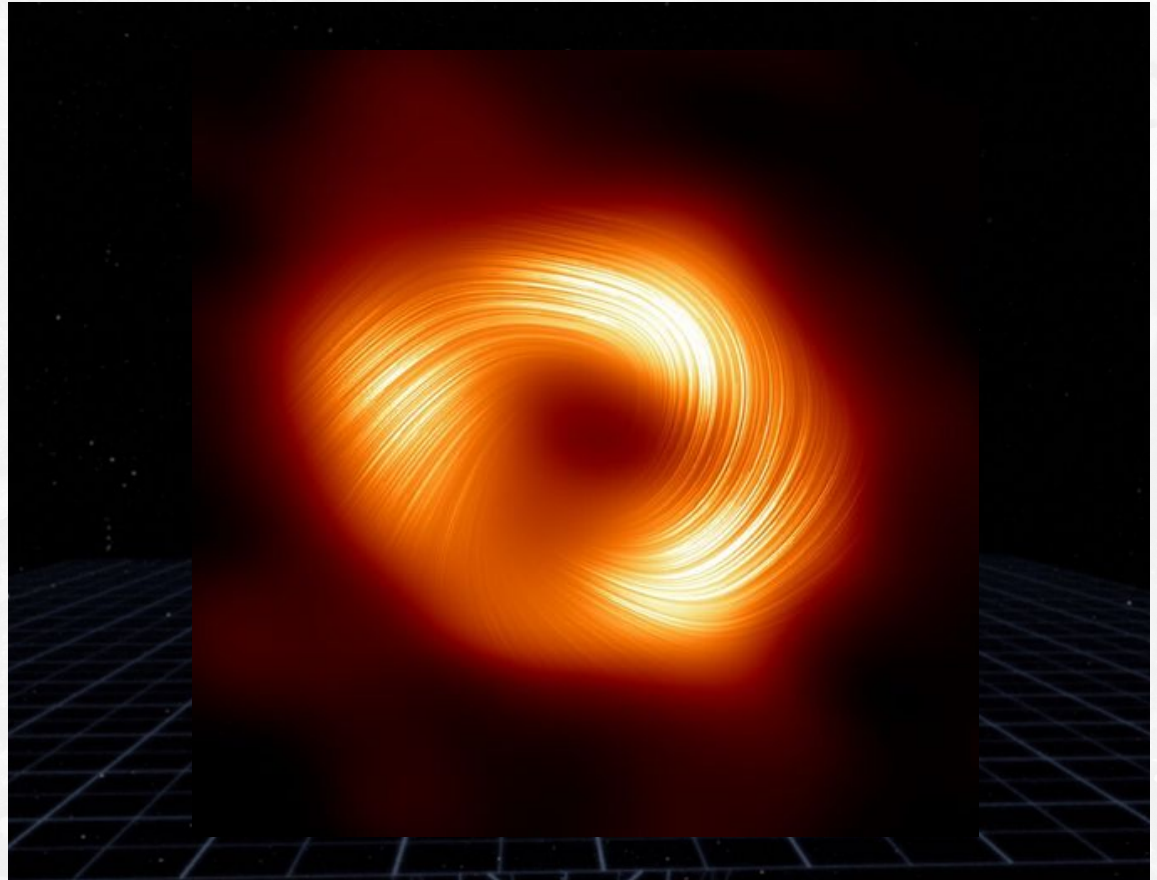
- But still impossible to disentangle with current resolution



Credit: Center for Astrophysics | Harvard & Smithsonian

Secondary image: the photon ring

- But still impossible to disentangle with current resolution

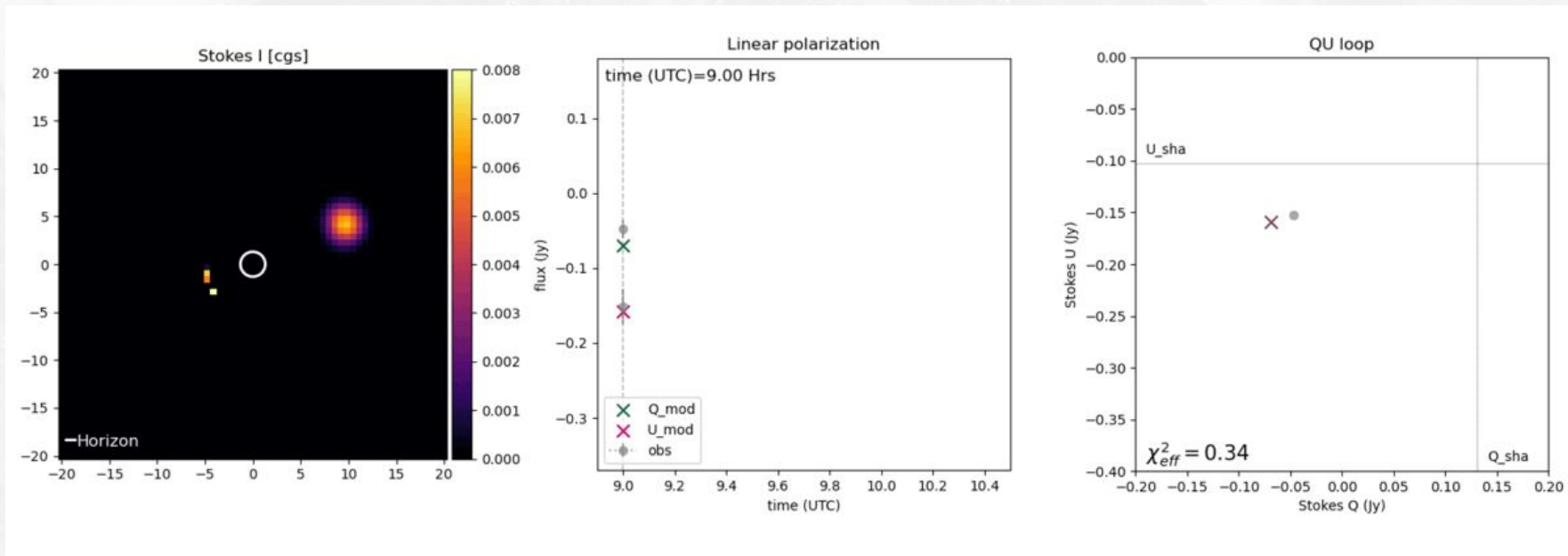
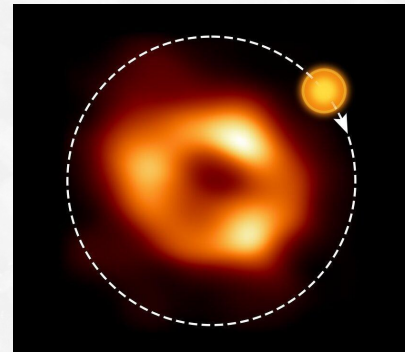


Credit: Center for Astrophysics | Harvard & Smithsonian

Why hotspot?

- ALMA fitting, Gravity fitting, Flares, Astrometry, Polarimetry, (VLBI?)
- Radio, NIR, Xray..

Credit:
EHT Collaboration, ESO/M. Kornmesser
(Acknowledgment: M. Wielgus)

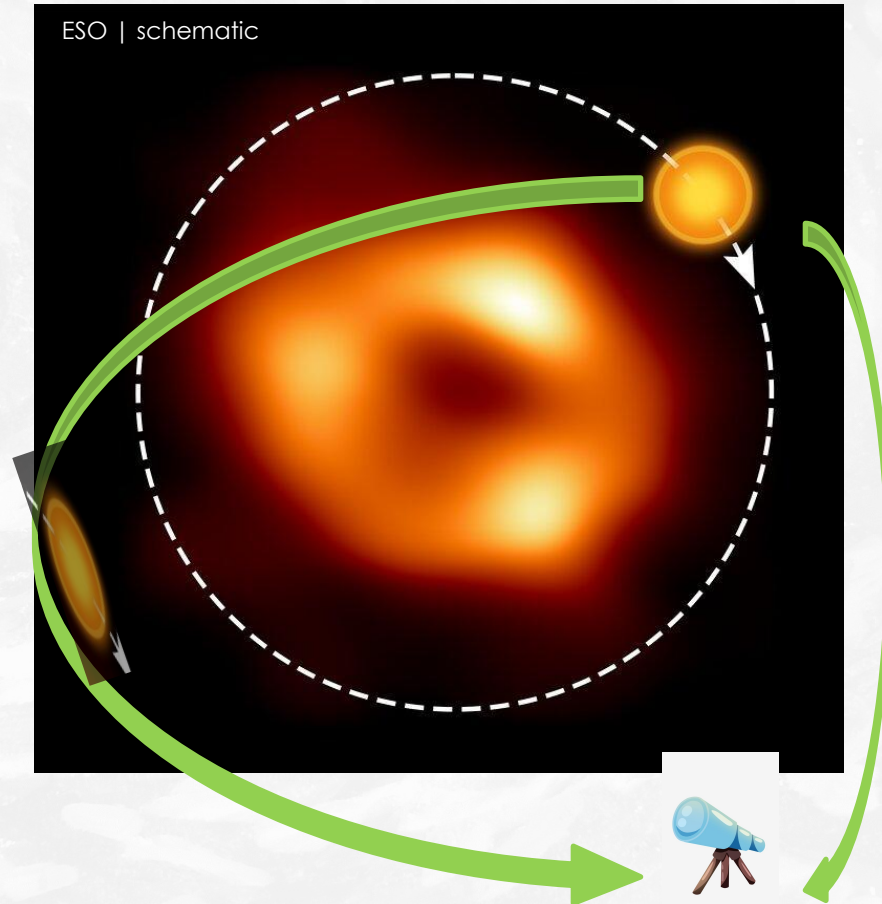


Yfantis+ 24a,b

Spin estimation from hot-spot observations

Yfantis+26

TWO SPOTS !!



Some solutions about it..

source and observer. By integrating along the trajectory, the geodesic equation (4) may be recast in integral form,⁴

$$I_r = G_\theta, \quad (7a)$$

$$\Delta\phi := \phi_o - \phi_s = I_\phi + \lambda G_\phi, \quad (7b)$$

$$\Delta t := t_o - t_s = I_t + a^2 G_t, \quad (7c)$$

where we define

$$I_r = \int_{r_s}^{r_o} \frac{dr}{\pm_r \sqrt{\mathcal{R}(r)}}, \quad (8a)$$

$$G_\theta = \int_{\theta_s}^{\theta_o} \frac{d\theta}{\pm_\theta \sqrt{\Theta(\theta)}}, \quad (8b)$$

$$I_\phi = \int_{r_s}^{r_o} \frac{a(2Mr - a\lambda)}{\pm_r \Delta(r) \sqrt{\mathcal{R}(r)}} dr, \quad (8c)$$

$$G_\phi = \int_{\theta_s}^{\theta_o} \frac{\csc^2 \theta}{\pm_\theta \sqrt{\Theta(\theta)}} d\theta, \quad (8d)$$

$$I_t = \int_{r_s}^{r_o} \frac{r^2 \Delta(r) + 2Mr(r^2 + a^2 - a\lambda)}{\pm_r \Delta(r) \sqrt{\mathcal{R}(r)}} dr, \quad (8e)$$

$$G_t = \int_{\theta_s}^{\theta_o} \frac{\cos^2 \theta}{\pm_\theta \sqrt{\Theta(\theta)}} d\theta. \quad (8f)$$

Here, the notation f indicates that these integrals are to be understood as path integrals along the photon trajectory, with the signs $\pm_r = \text{sign}(p^r)$ and $\pm_\theta = \text{sign}(p^\theta)$ switching at radial and angular turning points, respectively. In particular, all path integrals increase monotonically along the trajectory.

A. Angular integrals

The analysis of the angular integrals differs depending on the region of conserved quantity space. In this paper, unless otherwise specified, we will restrict to positive η ,

$$\eta > 0, \quad (9)$$

$$u_\pm = \triangle_\theta \pm \sqrt{\triangle_\theta^2 + \frac{\eta}{a^2}}, \quad \triangle_\theta = \frac{1}{2} \left(1 - \frac{\eta + \lambda^2}{a^2} \right). \quad (11)$$

To aid in the expression of the angular path integrals G_θ , G_ϕ , and G_t , we introduce the notation

$$F_i = F \left(\arcsin \left(\frac{\cos \theta_i}{\sqrt{u_\pm}} \right) \middle| \frac{u_\pm}{u_-} \right), \quad (12)$$

$$\Pi_i = \Pi \left(u_\pm; \arcsin \left(\frac{\cos \theta_i}{\sqrt{u_\pm}} \right) \middle| \frac{u_\pm}{u_-} \right), \quad (13)$$

$$E'_i = E' \left(\arcsin \left(\frac{\cos \theta_i}{\sqrt{u_\pm}} \right) \middle| \frac{u_\pm}{u_-} \right), \quad (14)$$

where $i \in \{s, o\}$ can be either source or observer. Here, $F(\varphi|k)$, $E(\varphi|k)$, and $\Pi(n; \varphi|k)$ respectively denote the incomplete elliptic integrals of the first, second, and third kind,⁵ while the prime denotes a derivative with respect to k , $E'(\varphi|k) := \partial_k E(\varphi|k) = [E(\varphi|k) - F(\varphi|k)]/(2k)$. These integrals vanish at the equator,

$$F_i = \Pi_i = E'_i = 0, \quad (\theta_i = 0) \quad (15)$$

and become complete at turning points,

$$F_i = \mp K, \quad \Pi_i = \mp \Pi, \quad E'_i = \mp E, \quad (\theta_i = \theta_\pm) \quad (16)$$

where our notation for the complete elliptic integrals is

$$K = K \left(\frac{u_+}{u_-} \right) = F \left(\frac{\pi}{2} \middle| \frac{u_+}{u_-} \right), \quad (17)$$

$$\Pi = \Pi \left(u_+ \middle| \frac{u_+}{u_-} \right) = \Pi \left(u_+; \frac{\pi}{2} \middle| \frac{u_+}{u_-} \right), \quad (18)$$

$$E' = E' \left(\frac{u_+}{u_-} \right) = E' \left(\frac{\pi}{2} \middle| \frac{u_+}{u_-} \right). \quad (19)$$

The $\eta > 0$ angular path integrals may be written in terms of these quantities and the number m of angular turning points encountered along the trajectory as [11, 26]

$$G_\theta = \frac{1}{a\sqrt{-u_-}} [2mK \pm_s F_s \mp_o F_o], \quad (20)$$

$$G_\phi = \frac{1}{a\sqrt{-u_-}} [2m\Pi \pm_s \Pi_s \mp_o \Pi_o], \quad (21)$$

$$G_t = -\frac{2u_+}{a\sqrt{-u_-}} [2mE' \pm_s E'_s \mp_o E'_o], \quad (22)$$

Finally, note that the integral for G_θ can be inverted to solve for θ_o or θ_s as a function of G_θ [11, 25, 26]. Since in this paper, we mainly fix the observer point (a telescope at infinity), we present θ_s in terms of θ_o and G_θ . This may be inferred from expressions for $\theta_o(G_\theta, \theta_s)$ by interchanging s and o , before sending $G_\theta \rightarrow -G_\theta$ to compensate.⁶ From Eq. (71) of Ref. [11] (noting that τ therein denotes G_θ , while ν_θ therein denotes $\pm_s$), we find

$$\frac{\cos \theta_s}{\sqrt{u_+}} = \text{sn} \left(F_o \pm_o \text{sign}(\eta) a \sqrt{-u_-} G_\theta \middle| \frac{u_+}{u_-} \right), \quad (25)$$

where $\text{sn}(\varphi|k)$ denotes the Jacobi elliptic sine function. This formula holds regardless of the sign of η [11, 26].

B. Radial integrals

In this paper, we will consider a distant observer at

$$r_o \rightarrow \infty. \quad (26)$$

Geodesics that reach this far observer have at most one radial turning point outside the horizon. Given a choice of conserved quantities (λ, η) , a simple way to test whether the ray has a turning point is to compute $r_4(\lambda, \eta)$ via Eq. (A8d) below. If r_4 is real and outside the horizon, then the ray has a turning point at radius r_4 ; otherwise, the ray never encounters a turning point.

For the rays with no turning point, the radial integrals I_r , I_ϕ , and I_t are single-valued functions of r_s , while for the rays with a turning point, these radial integrals must be double-valued in order to track whether or not the turning point has been reached. We will denote the number of turning points of a *photon* (portion of a ray) by $w \in \{0, 1\}$. The radial integral I_r may then be written

$$I_r = \int_{r_s}^{\infty} \frac{dr}{\sqrt{\mathcal{R}(r)}} + 2w \int_{r_4}^{r_s} \frac{dr}{\sqrt{\mathcal{R}(r)}}, \quad (27)$$

and likewise for I_ϕ and I_t with the appropriate integrands.⁷ We may relate w to the emission direction by

$$w = \begin{cases} 0 & p_r^* > 0, \\ 1 & p_r^* < 0 \text{ (and } r_+ < r_4 < r_s). \end{cases} \quad (28)$$

A ray reaching infinity (a photon) begins at the horizon (of the white hole) or from the associated radial integral I_r

$$I_r^{\text{total}} = \begin{cases} 2 \int_{r_4}^{\infty} \frac{dr}{\sqrt{\mathcal{R}(r)}} \\ \int_{r_+}^{\infty} \frac{dr}{\sqrt{\mathcal{R}(r)}} \end{cases}$$

where we remind the reader that $r_4(\lambda, \eta)$ is real and outside the horizon, and otherwise began at the horizon.

The full set of radial integrals reduced to elliptic form in Ref. [11] work in Refs. [24, 25]. The necessary computing Eqs. (27) and (29) are

As in Eq. (25) for $\theta_s(G_\theta)$, one formula for $r_s(I_r)$ [11, 25]. Eq. (25) a formula for r_o , and we may infer r_s described above Eq. (25), i.e., by sending $I_r \rightarrow -I_r$. Noting that and then sending $r_o \rightarrow \infty$, the emission

$$r_s = \frac{r_4 r_{31} - r_3 r_{41} \text{sn}^2 \left(\frac{1}{2} \sqrt{r_3} \right)}{r_{31} - r_{41} \text{sn}^2 \left(\frac{1}{2} \sqrt{r_{31}} \right)}$$

with

$$\mathcal{F}_o = F \left(\arcsin \sqrt{\frac{r_{31}}{r_{41}}} \middle| k \right)$$

Here, we introduced the notation

$$r_{ij} = r_i - r_j$$

with the roots $\{r_1, r_2, r_3, r_4\}$ given by This formula is contingent on the roots in the allowed range,

$$0 < I_r < I_r^{\text{to}}$$

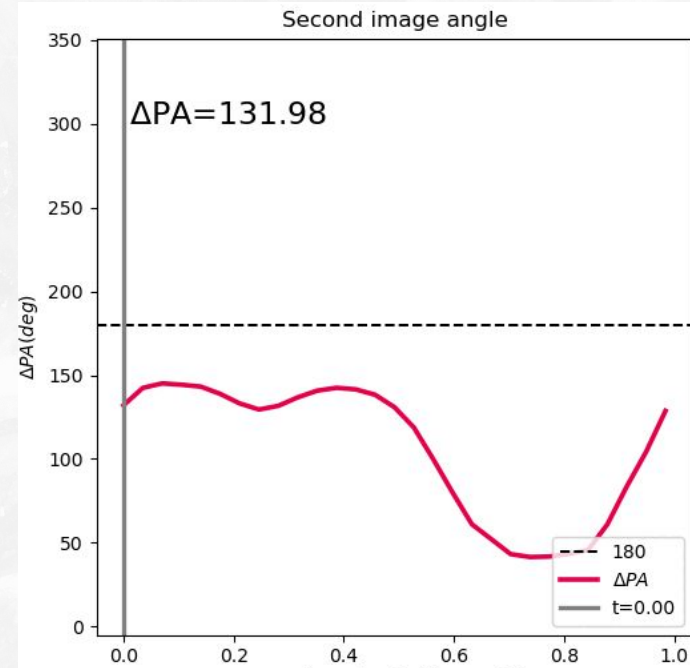
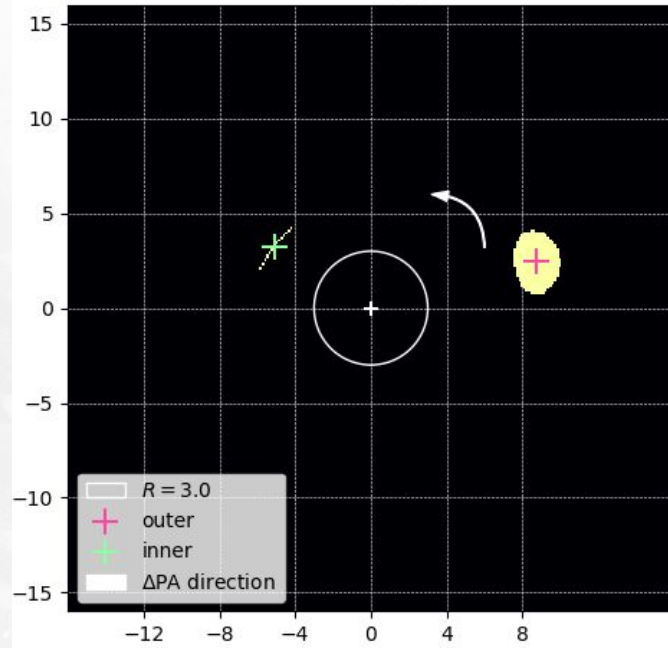
Provided that Eq. (33) is satisfied, the emission radius of a photon reaching the observer is a function of the conserved quantities (λ, η) . This formula (some of) the radial roots are co-

⁶The future-directed geodesic from source to observer is also a past-

Spin estimation from hot-spot observations

Yfantis, Palumbo, Moscibrodzka, 2026

distance (r_s) velocity (Ω) BH spin (a) viewing angle (i)
1000 combinations



Simulations in IPOLE

Spin estimation from hot-spot observations

Yfantis, Palumbo, Moscibrodzka, 2026

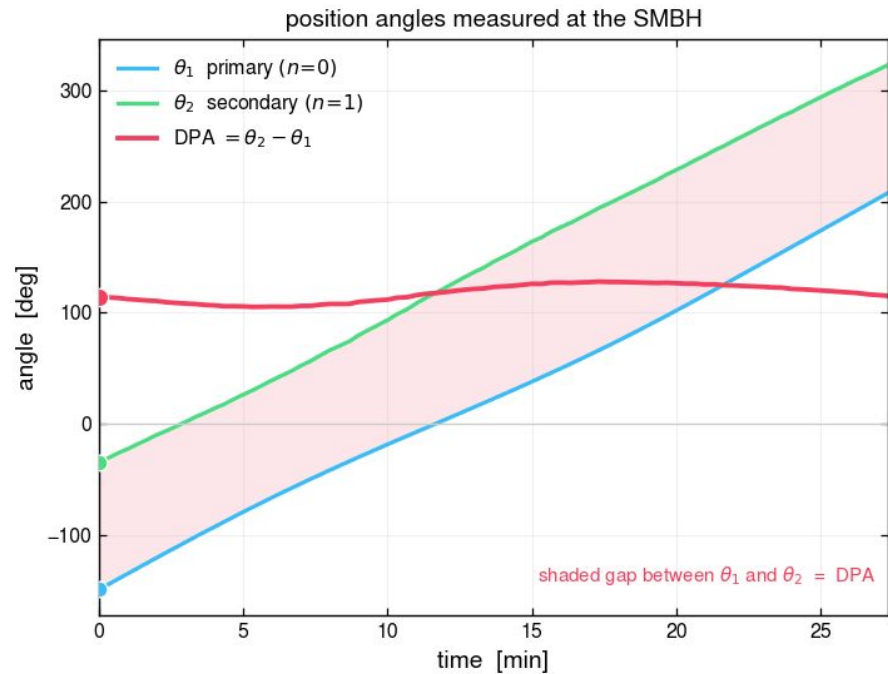
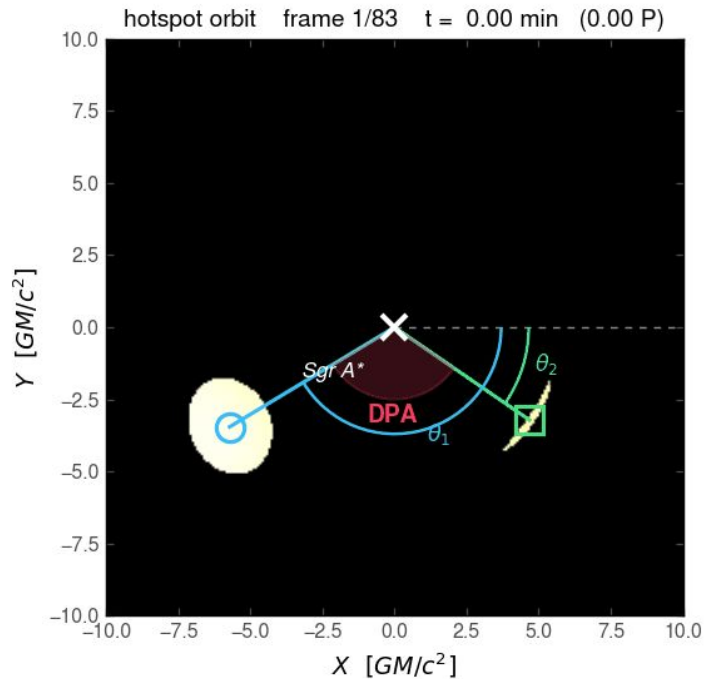
distance (r_s)

velocity (Ω)

BH spin (a)

viewing angle (i)

1000 combinations



Simulations in IPOLE

□ an empirical relation*

Yfantis+ 26

**BH
spin**

$$\Delta PA [deg] = 180 + a_* (40 - r_s) - \left(\frac{1408 - 233a_*}{r_s^2} + 35 \right) K_{\text{coeff}}$$

**Equatorial spots
i = 0 (or avg over full period)**

distance

speed

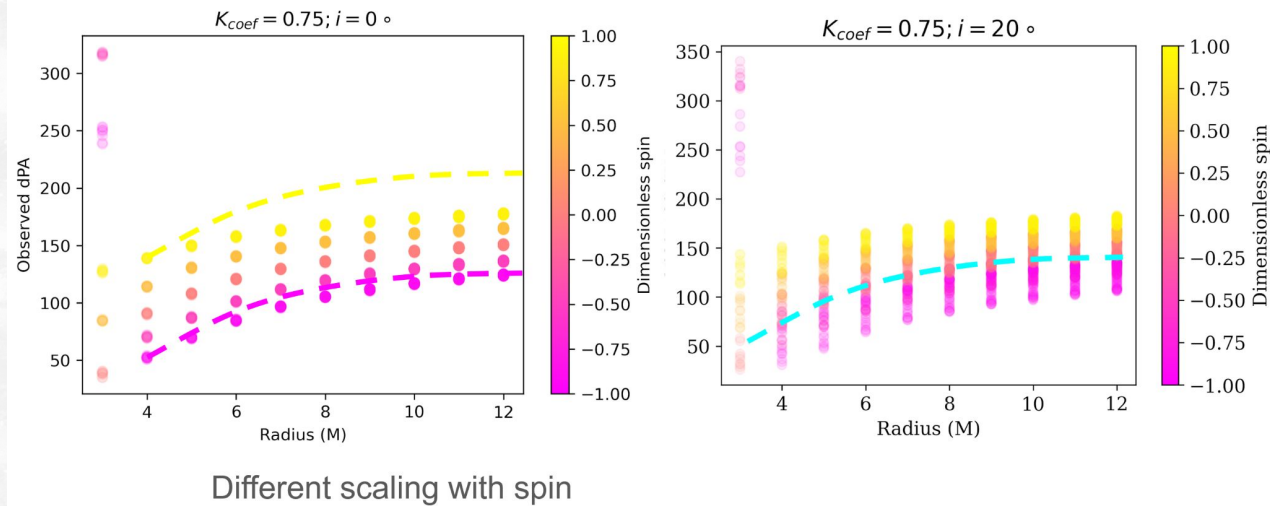
*Straight through pattern observations, trial-an-error and pyFit functions.

□ an empirical relation*

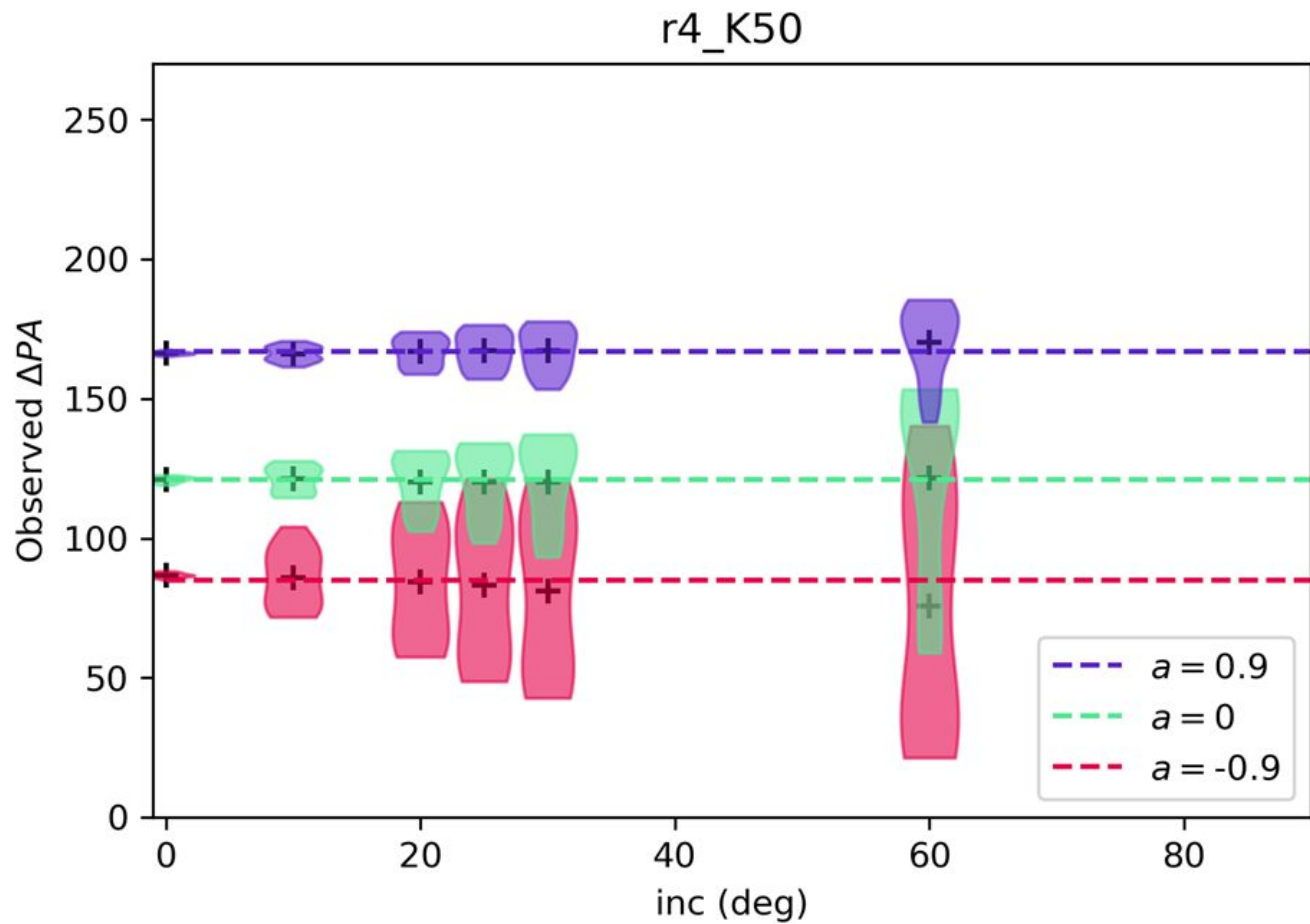
Yfantis+ 26

$$\Delta PA [deg] = 180 + a_* (40 - r_s) - \left(\frac{1408 - 233 a_*}{r_s^2} + 35 \right) K_{\text{coeff}}$$

***Straight through
pattern identification,
trial-and-error and
pyFit functions.**



Get rid of inclination



empirical relation

-comparison to analytic

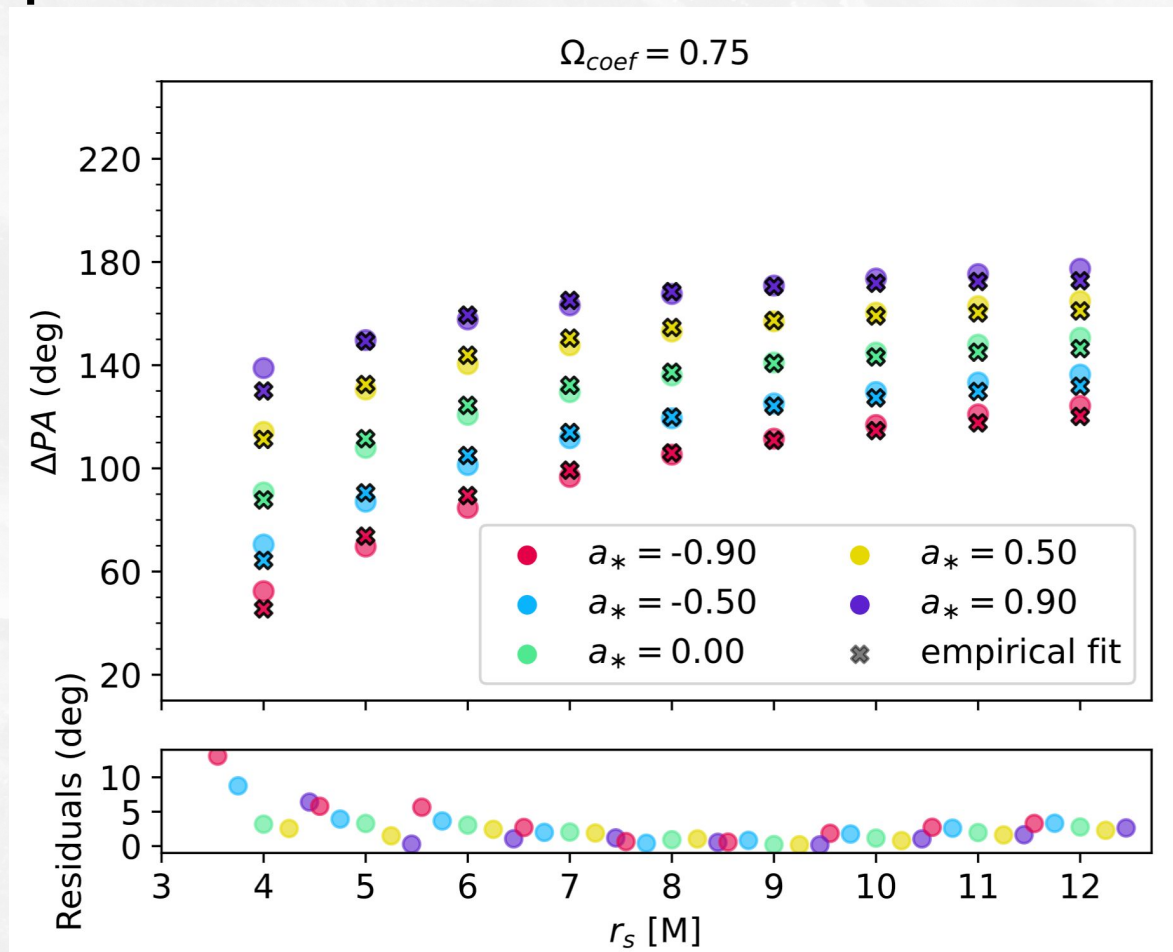
$$\Delta PA [deg] = 180 + a_* (40 - r_s) - \left(\frac{1408 - 233a_*}{r_s^2} + 35 \right) K_{\text{coeff}}$$

$$\delta PA_{0 \rightarrow 1} = \pi + \frac{2a}{\sqrt{\tilde{b}^2 - a^2}} \left(\frac{\tilde{r}_0 + M}{\tilde{r}_0 - M} \right) \times \\ \times K \left(\frac{\pi}{(\pi + a_* + 0.3)} \middle| \frac{a^2}{a^2 - \tilde{b}^2} \right).$$

$$\tilde{b} = \sqrt{\frac{\tilde{r}_0^3}{a_*^2} \left[\frac{4M\Delta(\tilde{r}_0)}{(\tilde{r}_0 - M)^2} - \tilde{r}_0 \right] + a_*^2}.$$

$$\tilde{r}_0 = M + 2 \sqrt{M^2 - \frac{a^2}{3}} \cos \left[\frac{1}{3} \arccos \left(\frac{1 - \frac{a^2}{M^2}}{\left(1 - \frac{a^2}{3M^2}\right)^{3/2}} \right) \right],$$

Performance of empiric relation



Next step

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**Astronomy
&
Astrophysics**

Lensing of hot spots in Kerr space-time An empirical relation for black hole spin estimation

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In review

A deep learning algorithm for black hole spin estimation using hot-spot secondary images

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*Authors have contributed equally in this work.



A deep learning algorithm for black hole spin estimation using hot-spot secondary images

A. I. Yfantis^{1*} and R. Al-Belmeisi^{2*}

Approximation by Superpositions of a Sigmoidal Function*

G. Cybenko†

Abstract. In this paper we demonstrate that finite linear combinations of compositions of a fixed, univariate function and a set of affine functionals can uniformly approximate any continuous function of n real variables with support in the unit hypercube; only mild conditions are imposed on the univariate function. Our results settle an open question about representability in the class of single hidden layer neural networks. In particular, we show that arbitrary decision regions can be arbitrarily well approximated by continuous feedforward neural networks with only a single internal, hidden layer and any continuous sigmoidal nonlinearity. The paper discusses approximation properties of other possible types of nonlinearities that might be implemented by artificial neural networks.

Key words. Neural networks, Approximation, Completeness.

- Deep Learning **can approximate any arbitrary function** [[G. Cybenko, 1989](#)]

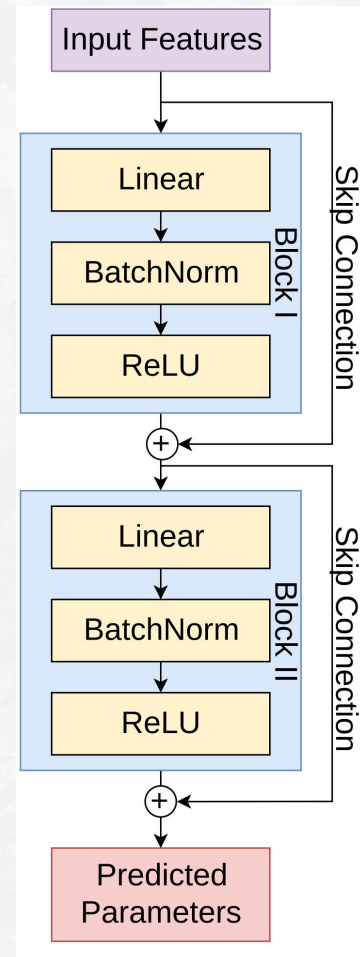
Our proposed framework

STIHOS: **S**pace**T**ime **I**nference via **H**ot**S**pot **O**bservations

*From the Greek στιχος = lyric

Deep Learning inference

Inputs / Observables	Inferred / Unknowns
r_{HS}	α
$\Delta PA(t)$	i
P	θ_z

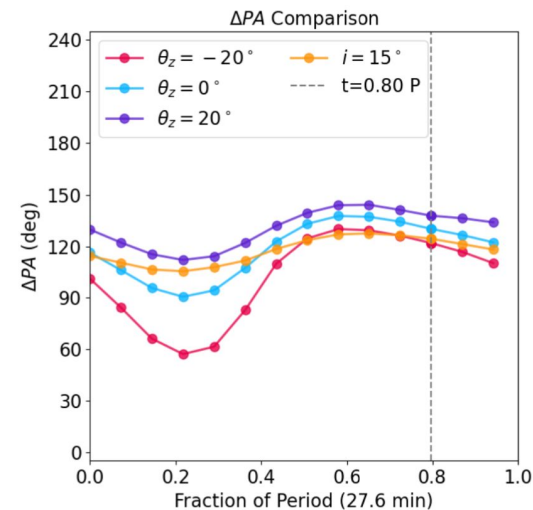
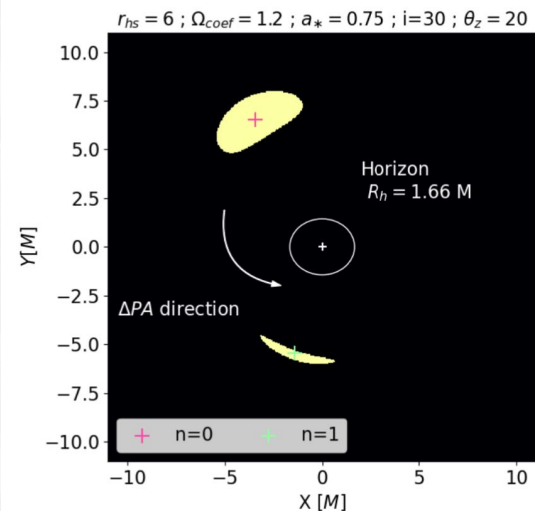


Explored Parameter Space

N = 100,000
(from 500
billion
combos)

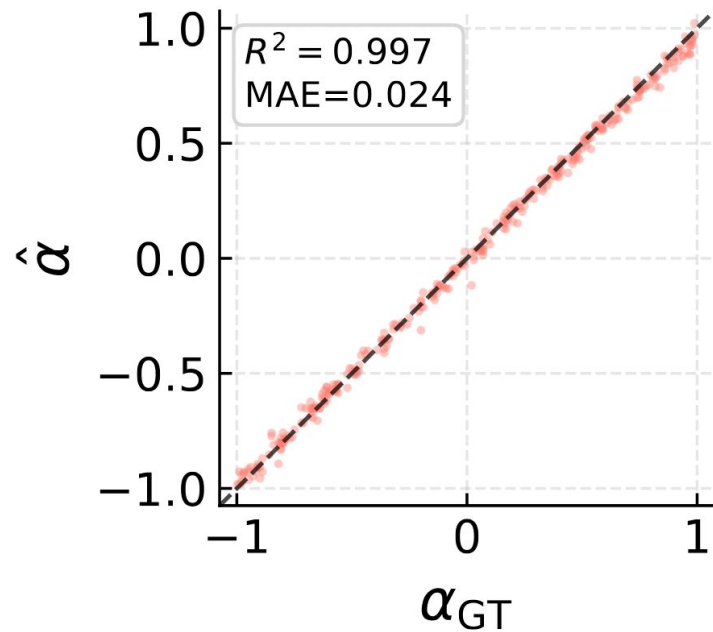
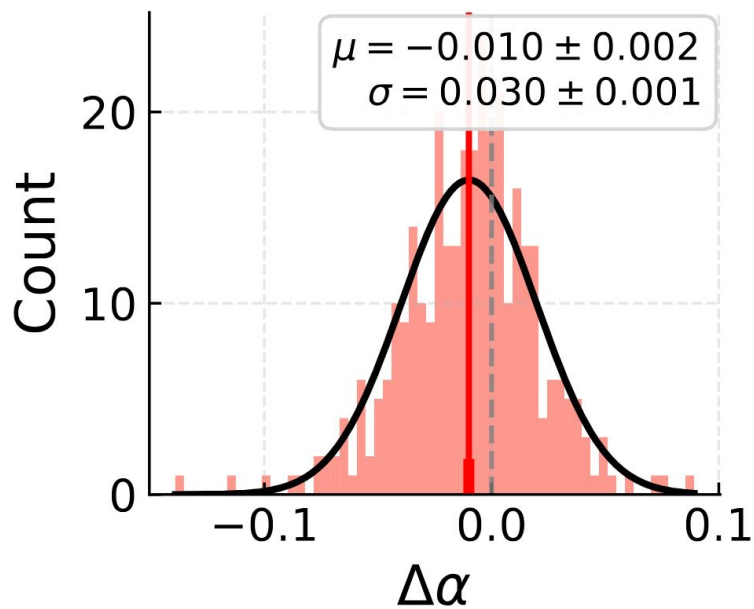
Parameter	Range	Cadence
a_*	$(-1, 1)$	0.01
r_{hs}	$[4, 12]$	0.1
Ω_{coef}	$[0.2, 1.5]$	0.1
i	$[0, 60]$	0.1
θ_z	$[-45, 45]$	1

Exp	Equatorial	Full orbit	Inputs	Predictions
I	✓	✗	$r, T, \Delta\text{PA}$	a_*
II	✓	✓	$r, T, \Delta\text{PA}(t)$	a_*, i
III	✓	✗	$r, T, \Delta\text{PA}(t)$	a_*, i
IV	✗	✓	$r, T, \Delta\text{PA}(t)$	a_*, i, θ_z
V	✗	✗	$r, T, \Delta\text{PA}(t)$	a_*, i, θ_z



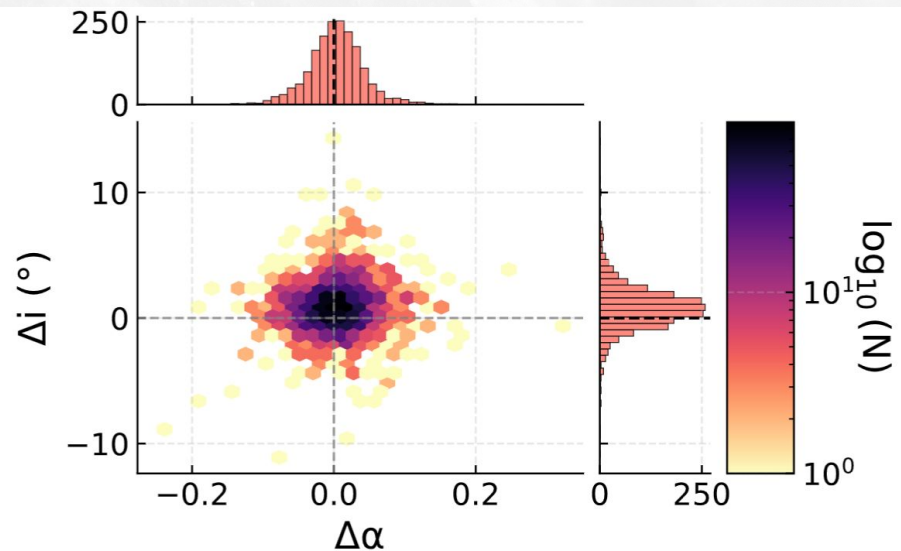
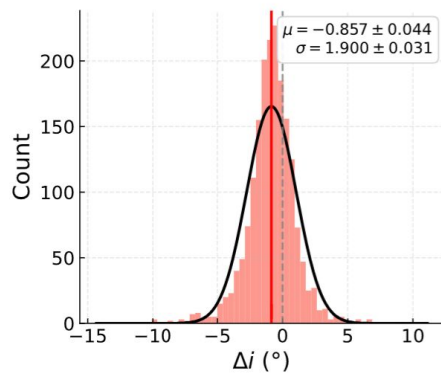
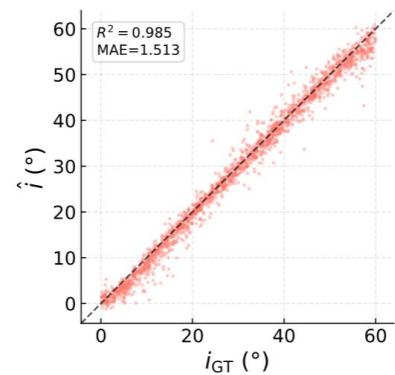
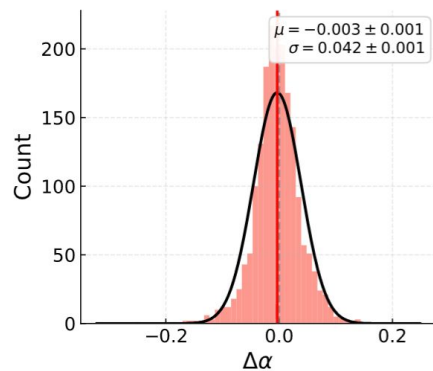
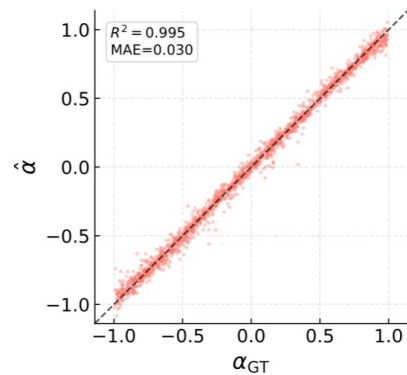
Results

Exp	Equatorial	Full orbit	Inputs	Predictions
I	✓	×	$r, T, \overline{\Delta PA}$	a_*



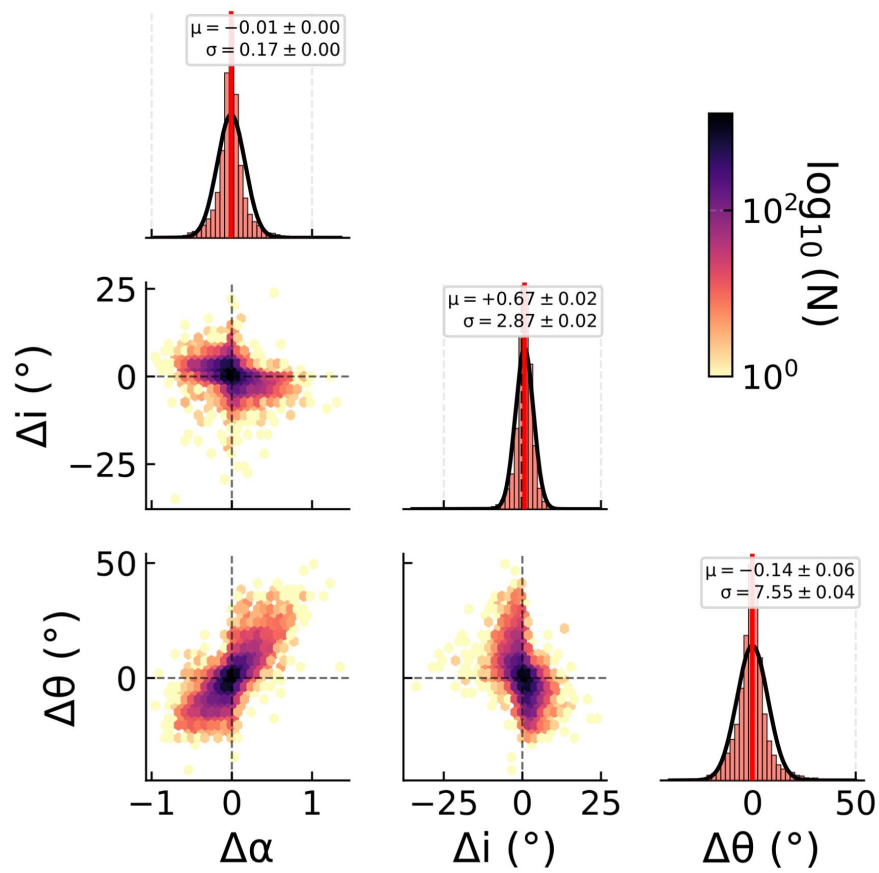
Results

Exp	Equatorial	Full orbit	Inputs	Predictions
II	✓	✓	$r, T, \Delta\text{PA}(t)$	a_*, i



Results

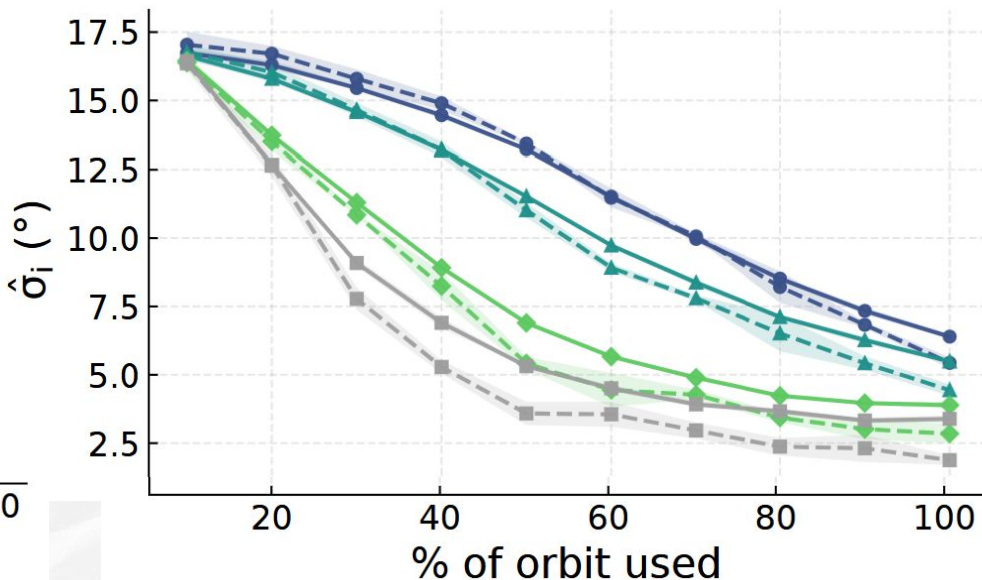
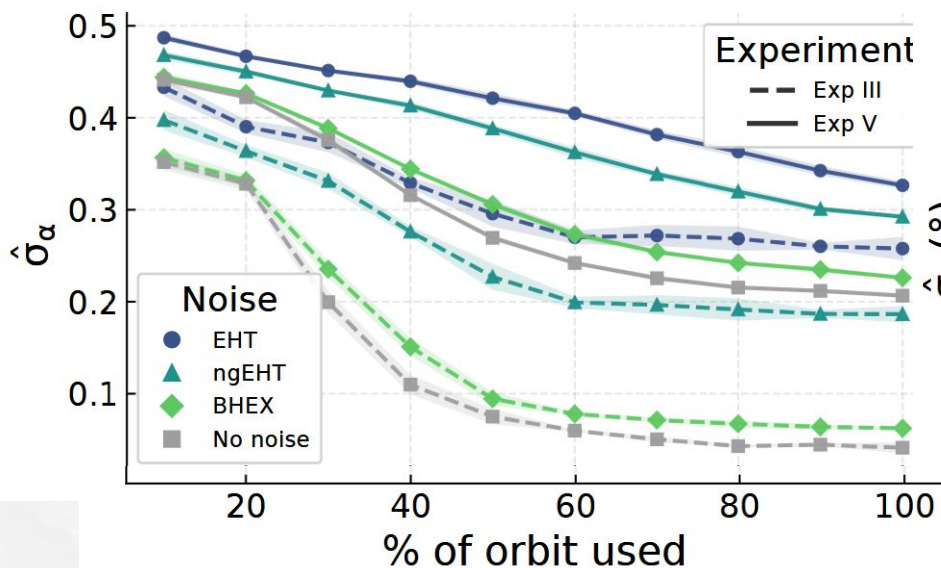
Exp	Equatorial	Full orbit	Inputs	Predictions
IV	×	✓	$r, T, \Delta\text{PA}(t)$	a_*, i, θ_z



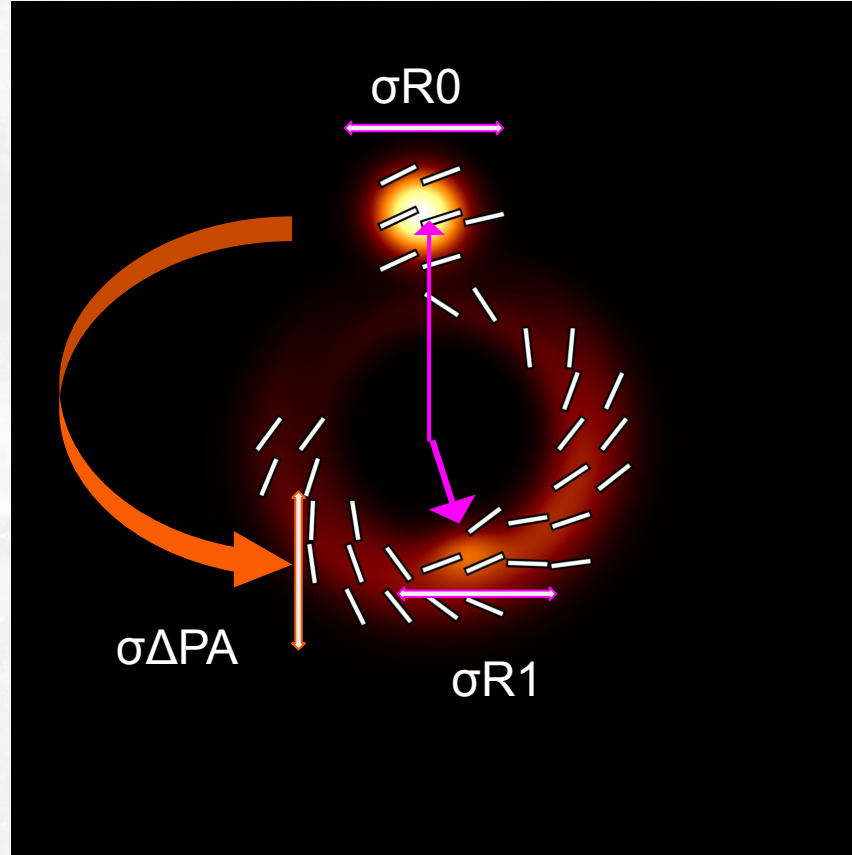
Results

Exp	Equatorial	Full orbit	Inputs	Predictions
III	✓	×	$r, T, \Delta PA(t)$	a_*, i

Part of orbits (from 10% to 100%)



How accurate constraints can one expect from realistic observations?



How accurate constraints can one expect from realistic observations?

Table 4. Virtual observations and assigned uncertainties.

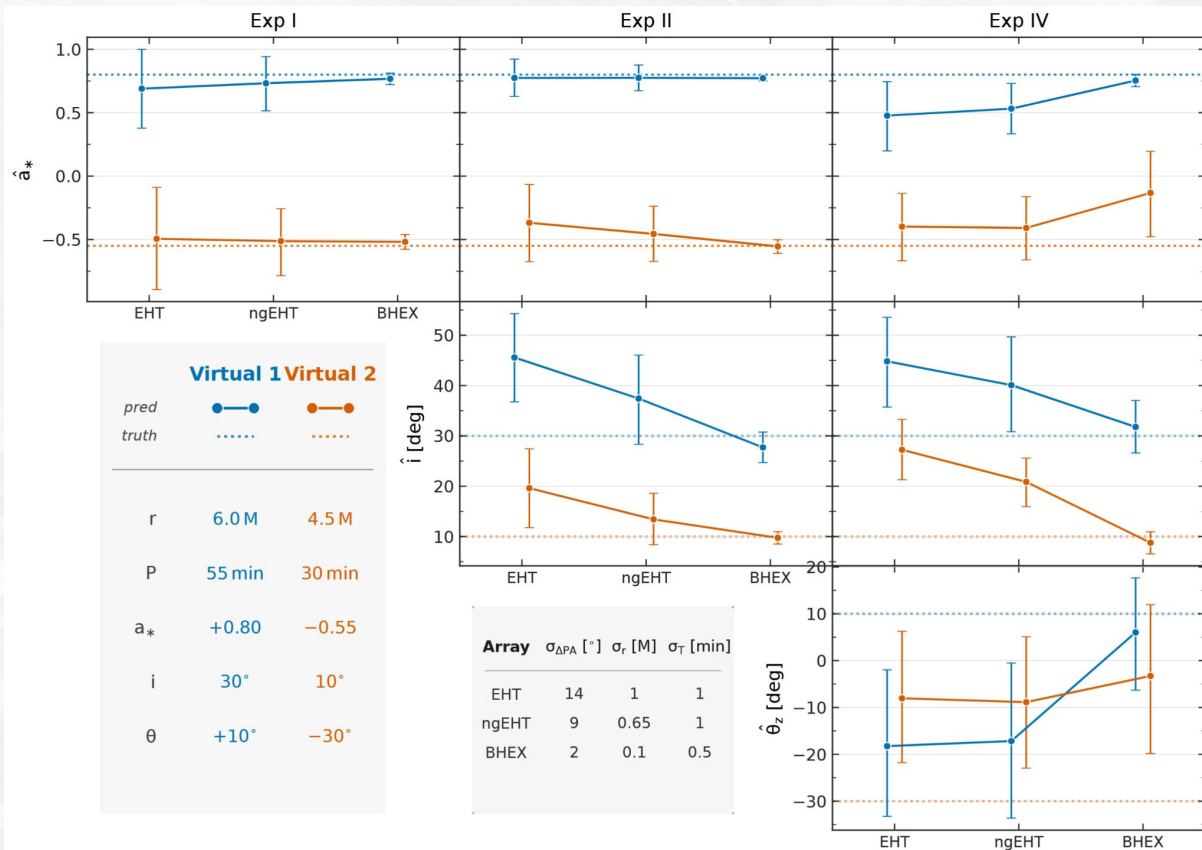
Instrument	$\sigma_{r_{\text{hs}}} [M]$	$\sigma_{\Delta PA} [^\circ]$	$\sigma_P [\text{min}]$
EHT	1.0	14	1
ngEHT	0.65	9	1
BHEX	0.1	2	0.5
Virtual data	$r_s [M]$	$\Delta PA [^\circ]$	$P [\text{min}]$
1	6.0	~ 170	55
2	4.5	~ 90	30

$$\Delta PA = \arctan\left(\frac{\Delta y}{\Delta x}\right), \quad \sigma_{\Delta PA} = \frac{2\sigma_{\Delta xy}}{L} [\text{rad}].$$

$$d = 51.8 \pm 2.3 \mu\text{as} \text{ (EHTC et al. 2022b,a)}$$

How accurate constraints can one expect from realistic observations?

Param.	Exp.	Virtual	EHT	ngEHT	BHEX
a_*	I	1	0.31	0.21	0.05
	I	2	0.44	0.31	0.08
	II	1	0.13	0.10	0.02
	II	2	0.28	0.20	0.05
	IV	1	0.25	0.19	0.04
	IV	2	0.27	0.24	0.3
$i[^\circ]$	II	1	9.3	8.9	2.3
	II	2	7.6	4.9	1.0
	IV	1	7.9	8.2	3.8
	IV	2	5.8	4.8	1.9
$\theta_z[^\circ]$	IV	1	14	14	8.6
	IV	2	13	12	11

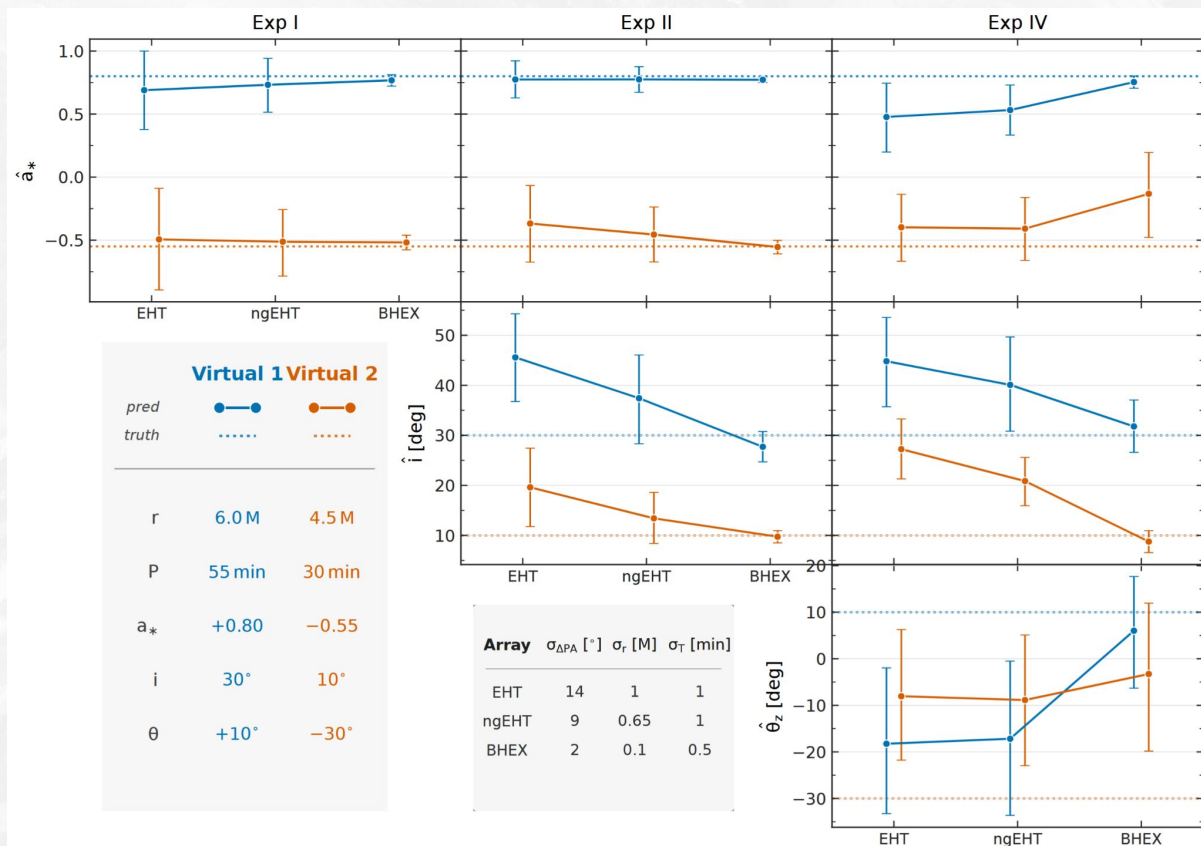


How accurate constraints can one expect from realistic observations?

Estimates can be improved from

Recurring observations

Prior info on inc, hotspot location (jet vs disk) ..



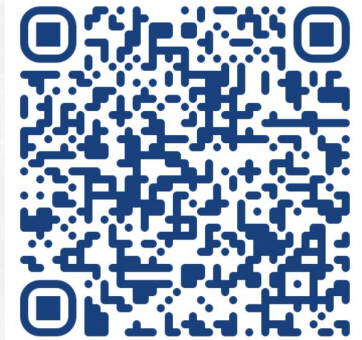
Observational feasibility

Or Looking forward

- Feature extraction: Identify hot spots
- Extract the observables ($r, P, \Delta PA$ curves)
- Study of secondary image (time delay, Intensity ratios, movement)
- Ratio of full-orbits
- Include light-curves if possible (e.g. ALMA)

STIHOS AVAILABLE AT:

<https://github.com/RamiAlBel/SMBH-HotSpot-Orbits/tree/main/data>



In the Future..





The Black Hole Explorer

The Black Hole Explorer (BHEX) will reveal fundamental properties of supermassive black holes and their relativistic jets by creating a telescope larger than the Earth. BHEX will be proposed as a NASA Small Explorer Mission in 2025/2027.

TRACING THE EDGE OF THE UNIVERSE

How do black holes look?

A black hole's photon ring is a narrow, bright ring of light produced by photons that have orbited outside the black hole before escaping its gravity. By resolving this ring for the first time, BHEX will precisely measure a black hole's spacetime properties, including its mass and spin.

How do black holes shine?

Some supermassive black holes launch powerful relativistic jets that act as cosmic particle accelerators and power their entire host galaxies. BHEX will reveal how black holes launch and power these jets by surveying a population of blazars, making moves like: correct spinning black holes to their jets.

How do black holes grow?

BHEX will target a diverse population of supermassive black holes in a variety of accretion states, measuring black holes across different stages of their life cycles.



- 34m antenna
- 100 + 300 GHz simultaneous dual-band observation
- 100 Gb/s downlink using laser communications
- 20,000 km orbital altitude
- 6 μas angular resolution via space-ground interferometry
- 2032 launch for a 2-year mission



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UNIVERSITY



BLACK HOLE
INITIATIVE



National Radio
Astronomy
Observatory



MIT
HAYSTACK
OBSERVATORY



NAOJ
National Astronomical
Observatory of Japan



NASA

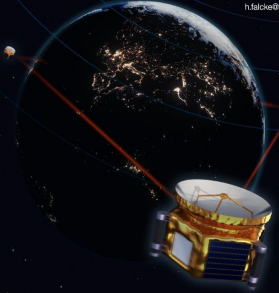
SHARP

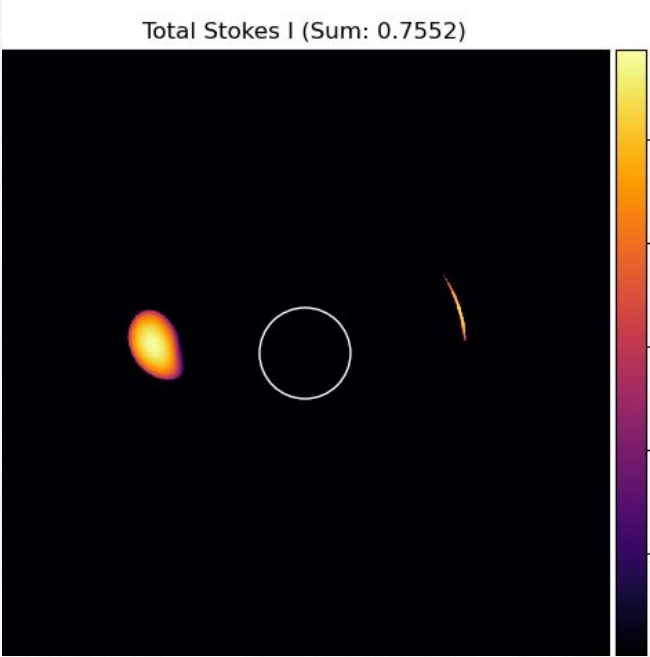
Space-based High-resolution Array
for Radio astronomy and Physics

M-class Mission Proposal

Lead proposer:
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Thank you!

Spin estimations from lensing of hot spot around SMBHs

From Empiric to Deep Learning

Aristomenis Yfantis

+R. Al-Belmpeisi, D. Palumbo, M. Moscibrodzka



Warsaw BH workshop
09/2026