DRD1 Collaboration Meeting

Capacitive-Sharing Pixel Micromegas Detector



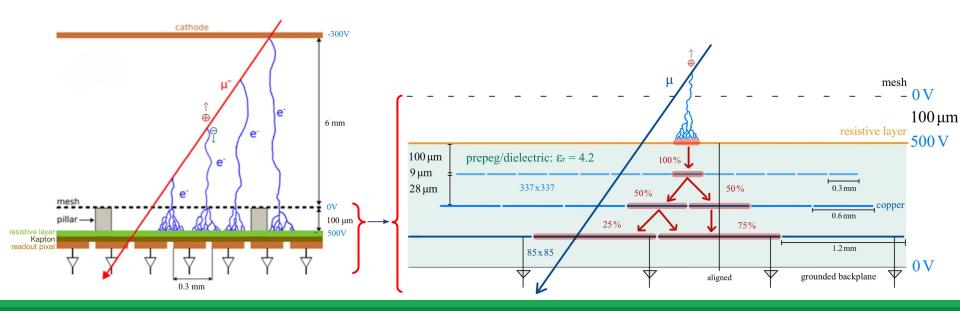
MAX MENZER

LMU: AG BIEBEL

8 October, 2025

Goal: Reduce the number of readout channels

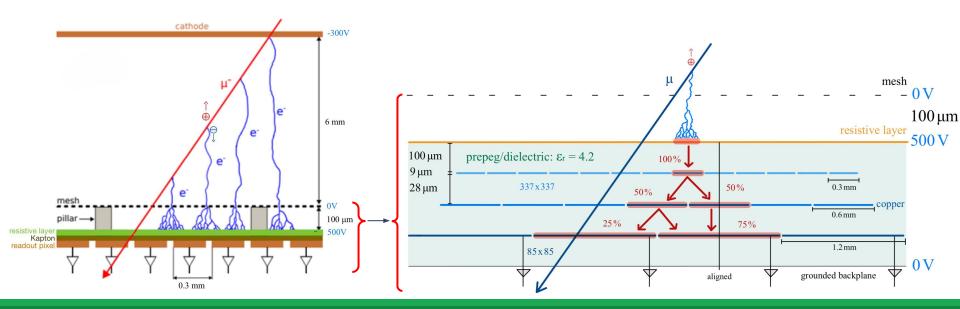
- Multiple layers of pixels that double in size
 → no electrical contact (dielectric)
- Readout only on the biggest pixels (readout pixels)



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3 layers: reduce readout pixels by factor $\sim 1/16$: $10 \text{ cm} \times 10 \text{ cm}$ detector $113500 \rightarrow 7225$



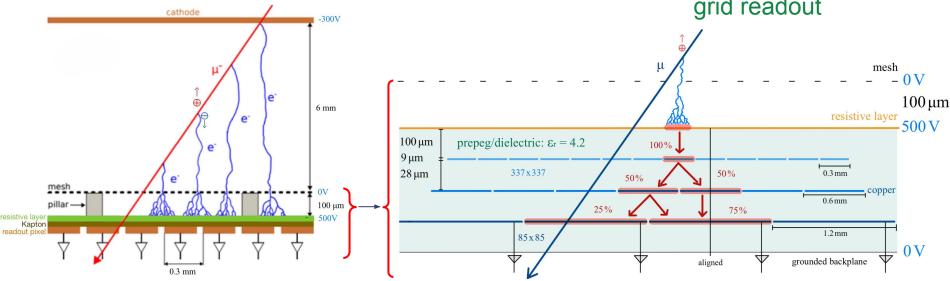
MAX MENZER

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3 layers: reduce readout pixels by factor ~ 1/16: 10 cm x 10 cm detector 113500 → 7225

Actual detector requires only 170 channels through a grid readout



Goal: Reduce the number of readout channels

Multiple layers of pixels that double in size
 → no electrical contact (dielectric)

100 μm

500V

- Readout only on the biggest pixels (readout pixels)
- Capacitive-Sharing allows the signal to travel downwards and split
- Unique readout for each position

25%

85 x 85

75%

aligned

Kapton

0V

1.2 mm

grounded backplane

Goal: Reduce the number of readout channels

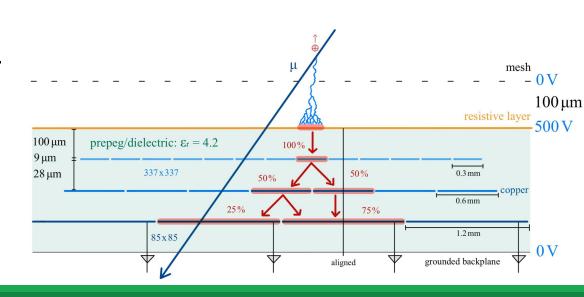
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weighted arithmetic mean

→ resolution in the order of the smallest pixels

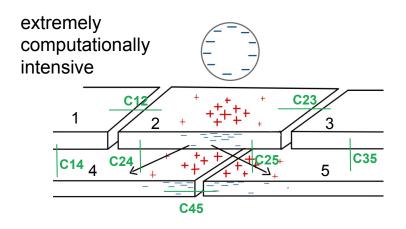
$$x_{rec} = \frac{\sum_{i=1}^{n} w_i x_i}{\sum_{i=1}^{n} w_i}$$

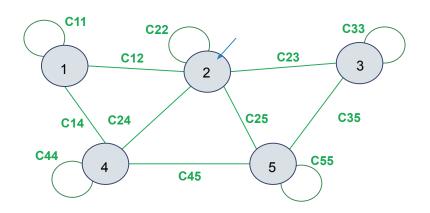
3 layers: reduce readout pixels by factor ~ 1/16: 10 cm x 10 cm detector 113 500 → 7225



Two Theories for Capacitive-Sharing

Equivalent Circuit	Ramo-Shockley Theorem
 capacitances can be calculated between each pair of pixels 	
 an equivalent circuit can replicate the detector 	
 a signal can be injected (pixel 2) and its propagation observed 	





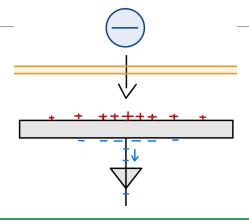
Two Theories for Capacitive-Sharing

Equivalent Circuit

- capacitances can be calculated between each pair of pixels
- an equivalent circuit can replicate the detector
- a signal can be injected (pixel 2) and its propagation observed

Ramo-Shockley Theorem

- a moving charge pushes charge carriers into the readout
- weighting field is a mathematical concept and a non real electric field
- ground all electrodes except electrode j, the one under consideration (1V)



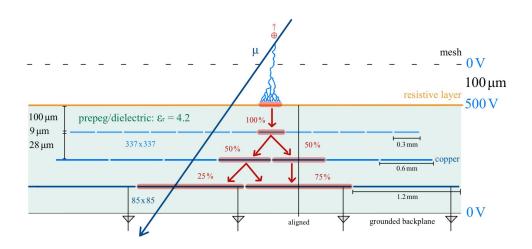
$$ij(t) = -q\mathbf{v}(t)\mathbf{E}(\mathbf{x}(t))w,j$$

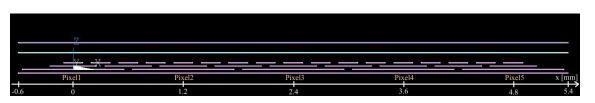
Ramo Shockley Theorem

(Notes Djunes Janssens [1])

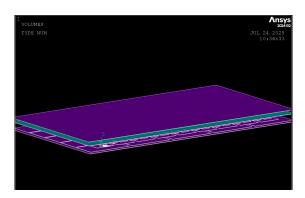
Setup for the ANSYS Simulation

Focus on a section of 5x5 readout pixels

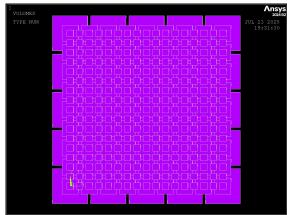




cross section with copper parts purple and dlc cyan



detector structure with mesh, dlc, pixels and backplane



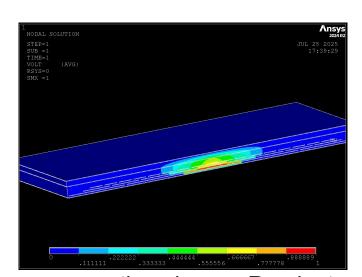
top view pixels only

$ij(t) = -q\mathbf{v}(t)\mathbf{E}(\mathbf{x}(t))w,j$

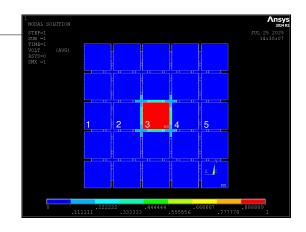
Weighting field for j=3

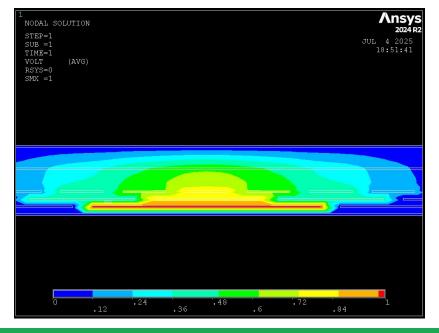
Electric potential with boundary conditions:

- Electrode/readout pixel j is set to unit potential (1V - red)
- All other electrodes are grounded
- The other pixel layers are floating



cross section along x: Readout pixels 1 to 5 visible

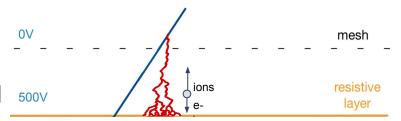




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Path integral Approach

Electrons drift to the dlc, ions to the mesh



Total induced charge is the integrated 500V
 Ramo-Shockley current

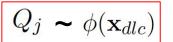
$$Q_j = Q_{j,e} + Q_{j,i} = \int_0^{t_e} i_{j,e}(t) dt + \int_0^{t_i} i_{j,i}(t) dt$$

Electron drift gives a path integral

$$Q_{j,e} = -q \int_0^{t_e} \mathbf{v}(t) \mathbf{E}(\mathbf{x}(t))_w dt = -q \int_0^{t_e} \frac{d\mathbf{x}(t)}{dt} \nabla \phi(\mathbf{x}(t))_w dt = -q \int_0^{t_e} \frac{d\phi(\mathbf{x}(t))_w}{dt} dt$$
$$= -e \left(\phi(\mathbf{x}_{dlc})_w - \phi(\mathbf{x}_0)_w\right)$$

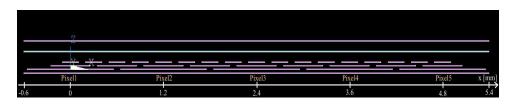
Add ion contribution to get total charge on readout pixel/electrode j

$$Q_j = -e \left(\phi(\mathbf{x}_{dlc})_w - \phi(\mathbf{x}_{mesh}) \right) \sim \phi(\mathbf{x}_{dlc})$$

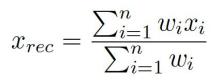


Path integral Approximation

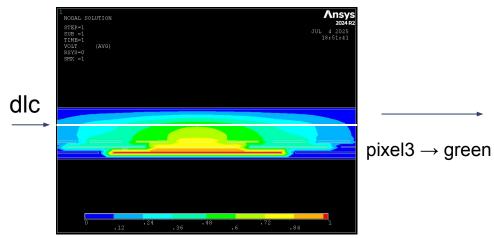
Use charge Qi as weighting wi for reconstruction in one dimension

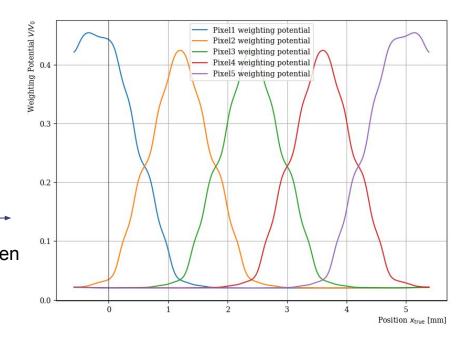


setup: cross section along x for reconstruction

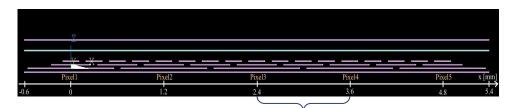


Project weighting potential at the dlc onto the y axis





Reconstruction

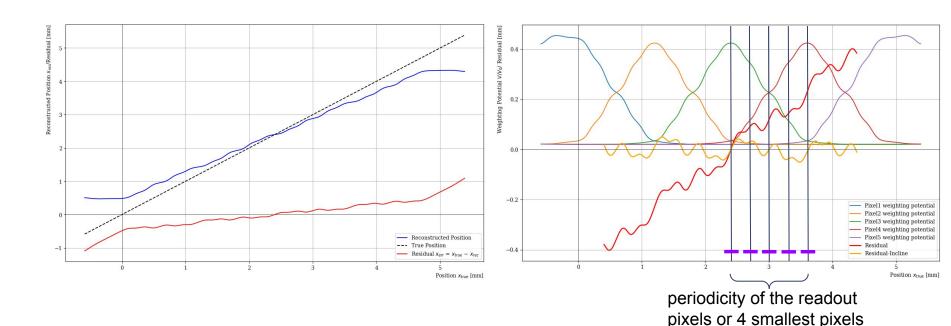


periodicity of the readout pixels or 4 smallest pixels

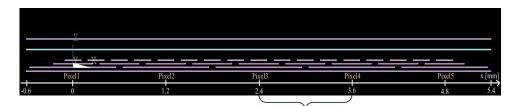
"Move" test charge from left to right:

Residual shows an inclination and oscillation

- Inclination results from limited size of the model (5x5)
- Oscillation from plateaus in the weighting potential



Reconstruction

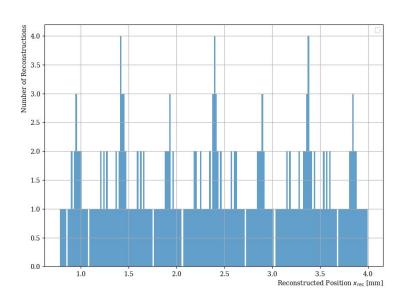


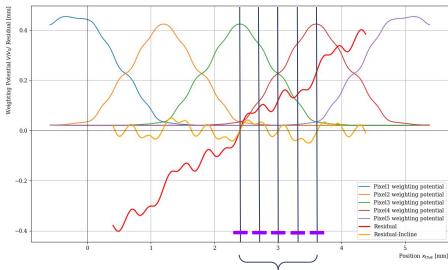
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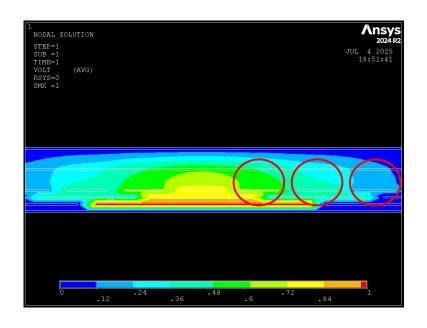


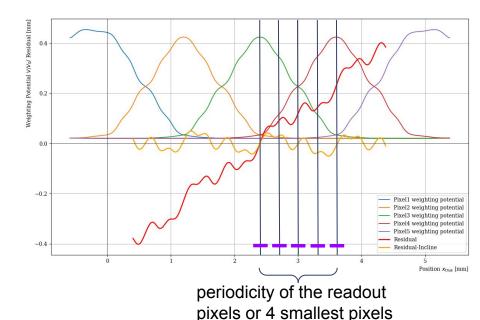


periodicity of the readout pixels or 4 smallest pixels

Oscillations of the Residual

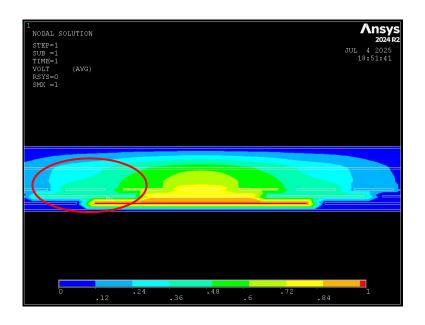
- weighting field E(x) (electric field) forces floating pixels to carry a constant potential
- expands E(x) but also propagates upwards (plateaus)

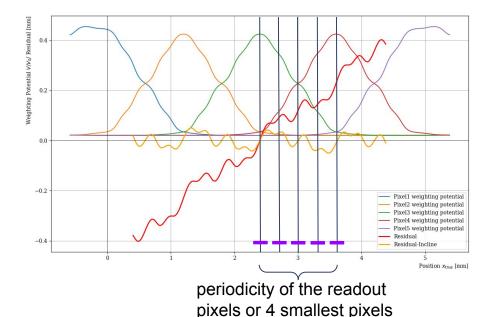




Oscillations of the Residual

- weighting field E(x) (electric field) forces floating pixels to carry a constant potential
- expands E(x) but also propagates upwards (plateaus)
- more prominent plateaus if pixel layers are aligned (substructure)
- substructure can be reduced by smaller pixel gap (notes Djunes)



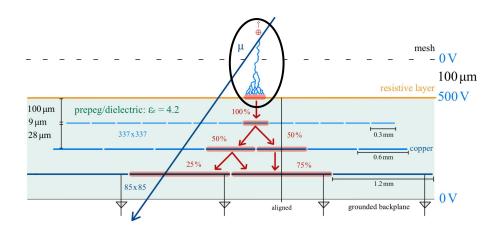


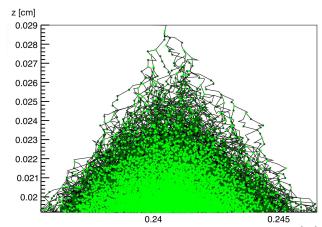
Microscopic Simulation - Garfield++

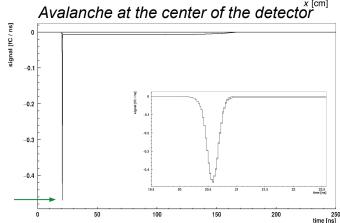
- Simulate detailed electron avalanche v(t)
- Import weighting field E(x(t)) from Ansys

Garfield can compute the current i_j(t) over time using the Ramo-Shockley Theorem

$$ij(t) = -q\mathbf{v}(t)\mathbf{E}(\mathbf{x}(t))w,j$$







Current from e-avalanche. Intense and short electron pulse with long ion tail

Microscopic Simulation - Garfield++

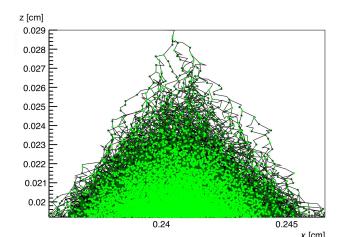
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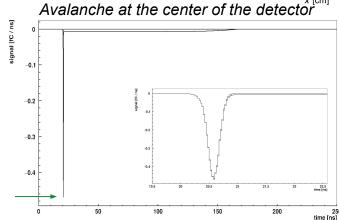
Garfield can compute the current i_j(t) over time using the Ramo-Shockley Theorem

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Maximum current ij(t) is used as weighting wi

$$x_{rec} = \frac{\sum_{i=1}^{n} w_i x_i}{\sum_{i=1}^{n} w_i}$$



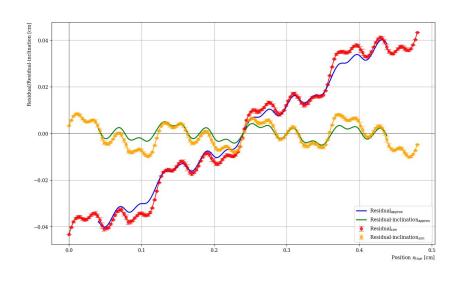


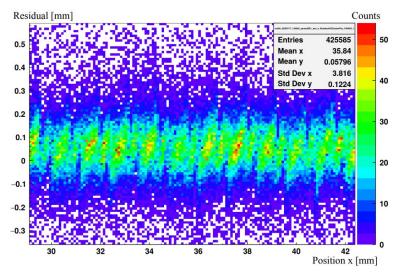
Current from e-avalanche. Intense and short electron pulse with long ion tail

Comparison and Measurement

Path integral and avalanche simulation exhibit the same behaviour as expected

Measurements H4 beamtime (CERN SPS 2025) show the same oscillations with substructure

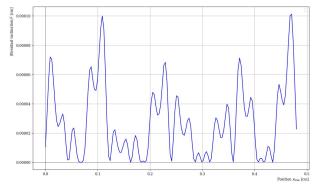




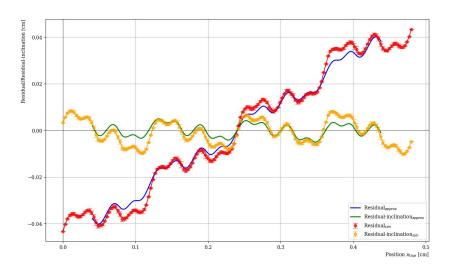
Comparison and Measurement

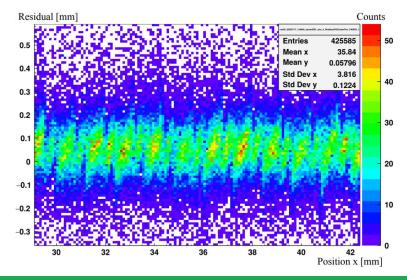
Path integral and avalanche simulation exhibit the same behaviour as expected

Measurements H4 beamtime (CERN SPS 2025) show the same oscillations with substructure The resolution is comparable to a single readout of the smallest pixels



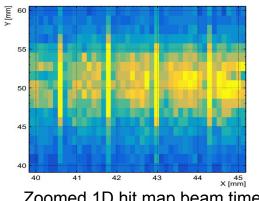
Estimated resolution using the Garfield++ simulation ~55µm



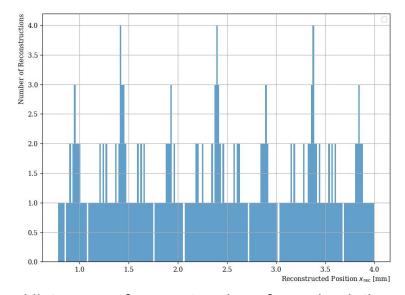


Reconstruction

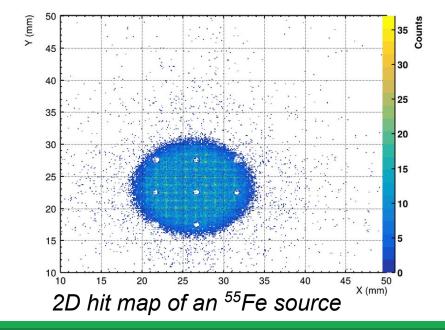
- Histogram of reconstructions and hit map show maxima in the center of the pixels
- Derivative of the residual also peaks there
- Simulation shows a slight deviation from the 1.2mm pitch (presumably limited model size)



Zoomed 1D hit map beam time

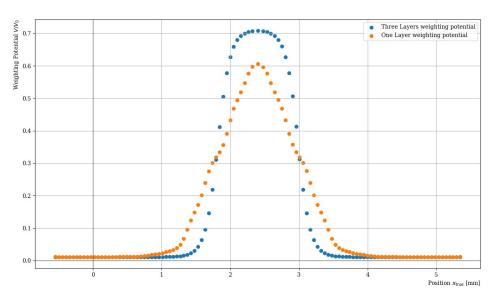


Histogram of reconstructions from simulation

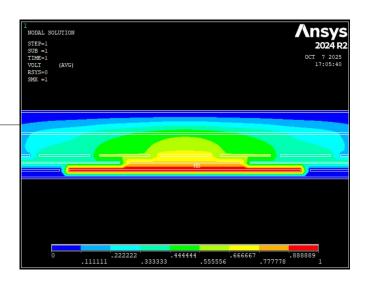


Loss of Signal Intensity

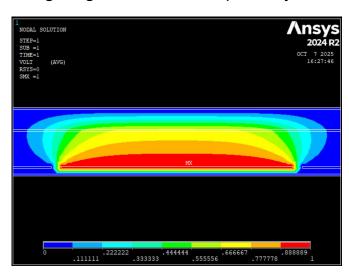
Pixel layers stretch and linearise the potential so we can use the weighted mean for reconstruction. Each layer stretches it further but also damps the signal more. Small plateaus and loss in intensity are the prize.



Projection of weighting field at the dlc



Weighting field with three pixel layers



weighting field without capacitive-sharing

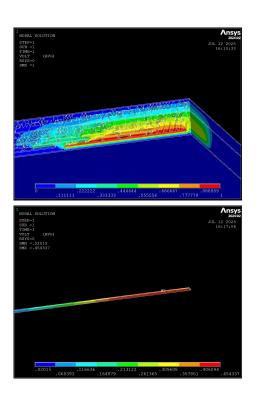
Backup

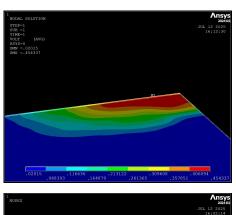
References for Capacitive Sharing

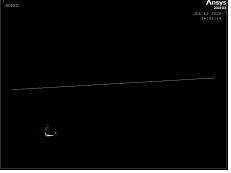
- [1] Unpublished notes and supporting ideas from Djunes Janssens (Ramo-Shockley Theorem)
- Original idea <u>Performance of a resistive micro-well detector with capacitive-sharing strip anode readout ScienceDirect</u>
- Studies on capacitive sharing resistive Micromegas for Muon System at FCC-ee.pdf
- Bachelor Thesis Max Menzer 2025, supervised by Eshita Kumar bac25_mmenzer.pdf, Bachelor Theses - Elementary Particle Physics - LMU Munich

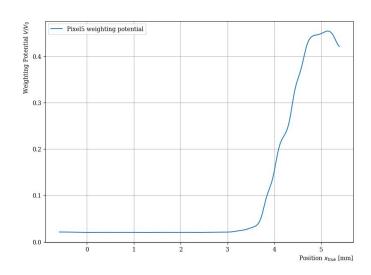
Path Integral Projection e.g. Pixel5

Weighting field in one dimension (x-axis) of pixel5 from Ansys







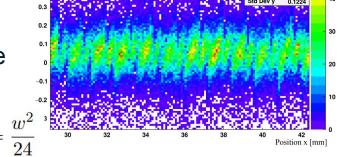


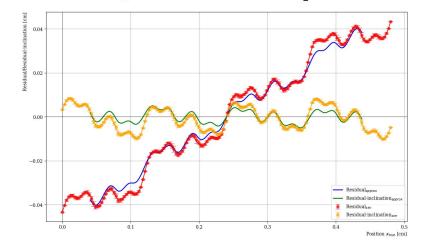
Comparison Residual

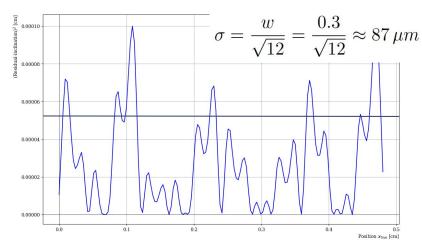
Using the residual (orange) to determine the resolution of the detector gives a limit of ~55µm

Shows that the resolution is even better than the classical readout without weighted mean

$$\sigma^2 = \frac{1}{w} \int_{-\frac{w}{2}}^{\frac{w}{2}} (x)^2 dx = \frac{1}{w} \left[\frac{x^3}{3} \right]_{-\frac{w}{2}}^{\frac{w}{2}} = \frac{1}{w} \frac{(w/2)^3}{3} - \frac{(-w/2)^3}{3} = \frac{w^2}{24}$$



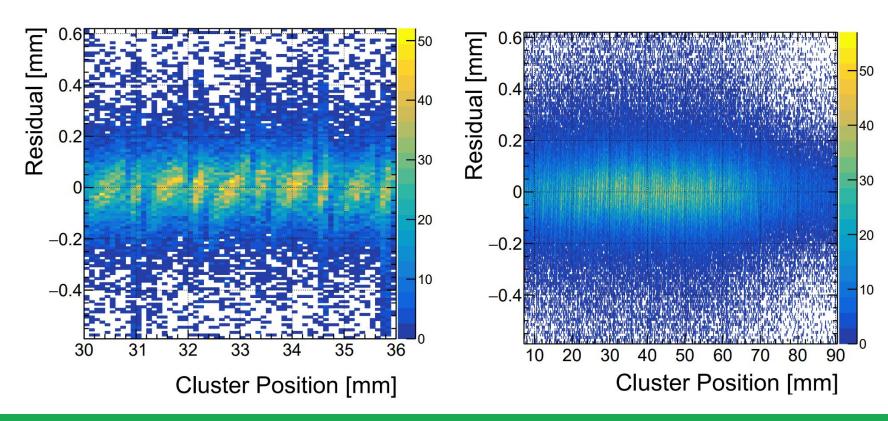


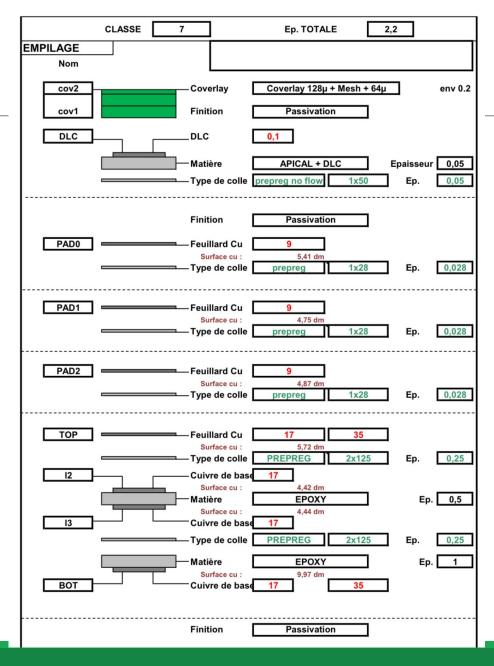


Squared residual (orange) for resolution

Additional Measurements H4 beamtime (CERN SPS25)

Measurements were carried out by Eshita Kumar and will be explained further tomorrow, Thu 09/10/25 11:45, DRD1



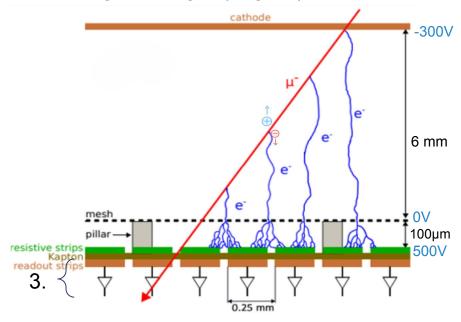


Stackup Three Layer Detector PADH used for measurements

MicroMegas (Micro-Mesh Gaseous Detector)

Goal: Position and Timing of particle passing through (e.g. μ)

- 1. Ionization and Drift region
- Charged particle ionizes gas
- E-field prevents recombination of e- and ions
 - → e- drift towards mesh
- 1. Amplification
 - Strong E-field between mesh and readout
 - e- avalanche
 - → enough charge to detect
- 1. Readout
- Transform "charge cloud" into signal
- Easiest: direct readout



<u>Problem:</u>

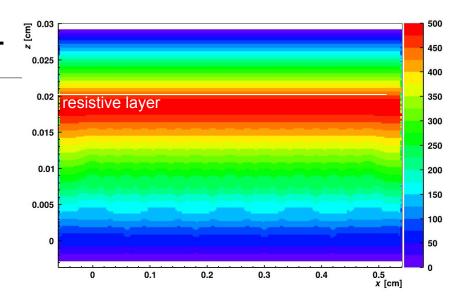
 many readout channels (electronics)

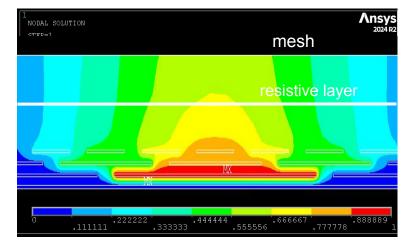
 $10x10cm \rightarrow \sim 150000$

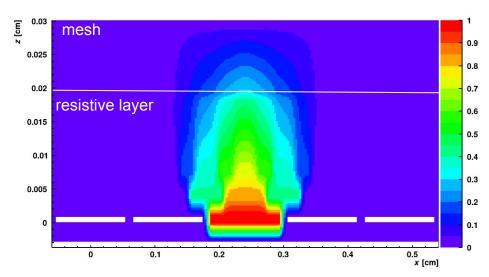
Import fields to Garfield++

Garfield - electric field (amplification)

weighting field (mathematical concept)



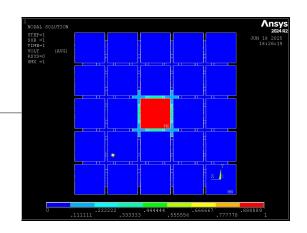




0.1 mm

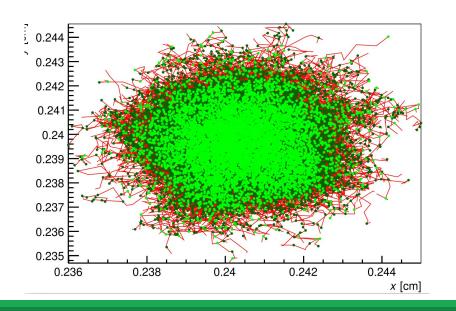
Electron avalanche

Centered above the center readout pixels

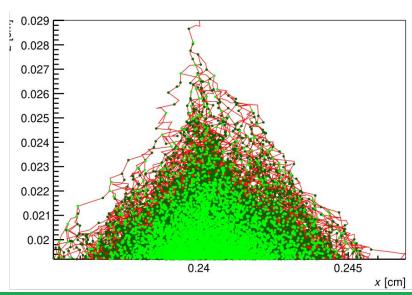


pixel size 1.2mm avalanche spread 0.1mm

top view (xy-plane



side view (xz-plane) cross section



Capacitive-Sharing Pixel Micromegas

Overview

- Micromegas Readout
- Capacitive-Sharing
- Explanation using the Ramo-Shockley Theorem
- Simulation using ANSYS
 - Path Integral
 - Microscopic Simulation (Garfield++)
- Comparison to Measurements

