

Energy flow and radiation luminosity in the simulations of neutron star Ultraluminous X-ray sources

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Neutron star Ultraluminous X-ray Sources (NS-ULXs)

- ULXs: Non-nuclear extra-galactic sources that emit X-rays at luminosities exceeding 10³⁹ erg s⁻¹, above the critical Eddington luminosity for a compact object with a mass less than 10 solar mass.
 - Super-Eddington accretion onto stellar mass objects (neutron stars and black holes),
 - Supper-Eddington radiation in the radii between magnetosphere and spherization radius,

 $L \simeq L_{Edd} [1 + \ln \dot{M} / \dot{M}_{Edd}],$ (Shakura & Sunyaev, 1973).

- Radiation pressure causes the outflow and resulting beaming (*King et al. 2001*).
- King. et.al, 2001: ULXs powered by extremely high accretion rate onto compact objects in high-mass X-ray binaries (HMXBs) in a transient stage.

Beaming emission

- KLK model (studied by King, Lasota & Kluźniak, 2017; King & Lasota, 2019, 2020):
 - ✓ The compact object emits its radiation within a fraction b of the unit sphere, causing the luminosity to be overestimated by a factor of 1/b
 - \checkmark when b<<1 then the inferred isotropic luminosity, $\rm L_{iso} \sim L/b$

 $L_{iso} = 4\pi d^2 F^r_{rad}$

 F^{r}_{rad} : radiation flux in optically thin cone-like region

d: distance between source and

detector



Simulation details

- General relativistic radiative magnetohydrodynamic (GRRMHD) code, KORAL.
- Conservation of mass ne energy-momentum tensor:

$$\nabla_{\mu}(\rho u^{\mu}) = 0$$
 & $\nabla_{\mu}(T^{\mu}{}_{\nu} + R^{\mu}{}_{\nu}) = 0$

- Neutron star: mass 1.4 M_{\odot} and radius $5r_g$ ($r_g = GM/c^2$) with dipole magnetic field
- Torus: Equilibrium torus with loop magnetic field with $\beta = (P_{gas} + P_{rad})/P_{mag} = 10$
- Resolution: $N_r x N_{\theta} x N_{\varphi} = 512 x 510 x 1$ with logarithmic spacing in r direction
- Schwarzschild metric
- Boundary condition: energy reflective surface for the neutron star with albedo 0.75

Initial setup



Models	B [G]	[L _{Edd} c⁻²]	M _{out} [L _{Edd} C ^{−2}]	L/L _{Edd}	b _{min}	L _{iso} /L _{Edd}	R _A /r _g
	1 × 10 ¹⁰	257	325	2.10	0.016	115	5.25
	2 × 10 ¹⁰	320	317	1.97	0.020	100	6.85
	3 × 10 ¹⁰	345	250	1.90	0.023	90	10.15
	5 × 10 ¹⁰	355	150	1.55	0.026	68	13.30
	7 × 10 ¹⁰	430	190	1.50	0.045	45	15.85
	1 × 10 ¹¹	490	80	1.49	0.083	39	17.39

3 × 10 ¹⁰	144	30	1.14	0.050	40	11.85
3 × 10 ¹⁰	1000	3600	2.5	0.010	185	8.10



Time-averaged data





Different regions

Be =

$$Be = -\frac{T^{t}_{t} + R^{t}_{t} + \rho u^{t}}{\rho u^{t}},$$

$$T^{t}_{t} \text{ and } R^{t}_{t} \text{ are MHD and radiation energy densities, } \rho \text{ is the mass-density and u^{t} is time component of 4-velocity}$$

$$0.01 - 45^{\circ}$$

$$0.01 - 45^{\circ}$$

$$0.01 - 45^{\circ}$$

$$0.00 - 45^{\circ}$$

$$0.01 - 45^{\circ}$$

$$0.00 -$$

0.02

$$\tau_{\rm r}(r) = \int_{r}^{r_{\rm out}} \rho \kappa_{\rm es} \sqrt{g_{\rm rr}} dr,$$

whit r_{out} the outer boundary of simulations, κ es = 0.34 cm g⁻¹, the electron scattering opacity for solar composition

0.02

0.01

0.00

-0.01

-0.02

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 $2\pi r \sin($

0

60°

75°

Radiation and kinatic efficiency

$$L_{\rm r} = -2\pi \int R^r{}_t \sqrt{-g} d\theta|_{\tau_{\rm r}<1}.$$

$$L_{\rm KE} = -2\pi \int (u_t + \sqrt{-g_{tt}})\rho u^r \sqrt{-g} d\theta \bigg|_{Be>0, u^r>0}$$

 $\eta \,{=}\, L/(\dot{M}c^2)$

Weaker magnetic field

- more powerful outflows
- higher radiation efficiency

Conclusions

- The simulation with a weak dipole (10^{10} G) compared to strong dipole (10^{11} G) :
 - ✓ More powerful outflow
 - Higher radiation efficiency
 - Higher apparent luminosity
- 10^{10} G =>> the apparent luminosity about 120 Eddington luminosity.
- 10^{11} G =>> the apparent luminosity about 40 Eddington luminosity.
- Increasing accretion rate in simulation with a fixed magnetic dipole strength:
 - More powerful outflow
 - Iower radiation efficiency
 - Higher apparent luminosity

The simulation with dipole strength 10¹⁰ G and accretion rate of 1000 Eddington units gives the apparent luminosity of 250 Eddington luminosity.