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Evidence for 3XMM J185246.6+003317 as a massive magnetar with a low magnetic field

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Layers of electron captures in the crust of accreting neutron stars

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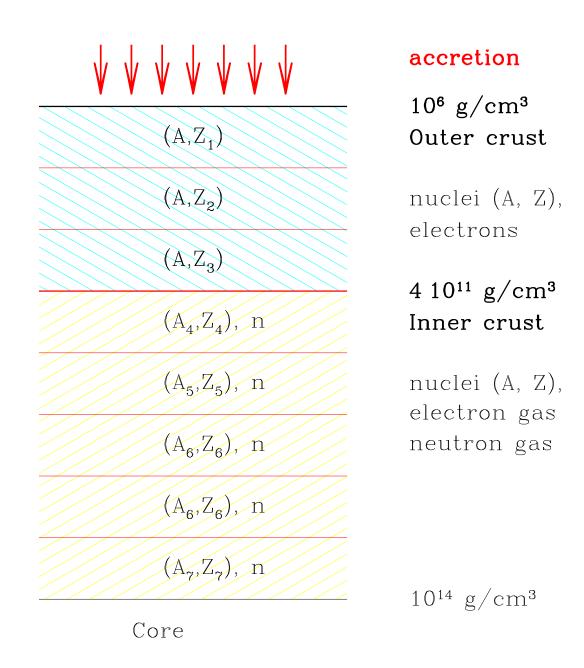
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Accretion vs. reactions

Classical approach	Reaction rate vs accretion
Infinite reaction rate	Reaction rate depends on experimental data, pressure and abundance of nuclei
Local equilibrium (t,P)	Matter not in equilibrium - reactions take place
Single nucleus approximation	In the reaction layer mixture of parent and grand-daughter nuclei (A,Z), (A,Z-2)
Reaction layer - infinitesimally thin	Time evolution of reaction layer (finite size)



Reaction Layer

$$(A, Z) + e^{-} \rightarrow (A, Z - 1) + \nu \text{ slow },$$

 $(A, Z - 1) + e^{-} \rightarrow (A, Z - 2) + \nu \text{ fast}$

Mixture of two nuclides (isobars)

Parent nuclei N_0 (A, Z). Grand-daughter nuclei N_2 (A, Z-2)

$$X = \frac{N_0}{N_0 + N_2}$$

 P_{th} threshold pressure - $\mu(P_{th},A,Z)=\mu(P_{th},A,Z-1)$ - corresponds to infinite reaction rate $\delta P=P-P_{th}>0$

The accretion timescale:

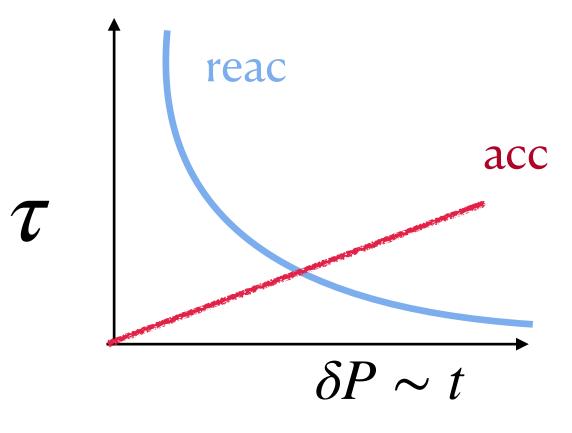
$$\tau_{acc}(\delta P) = \frac{4\pi R^4 \delta P}{GM\dot{M}}$$

- increasing function of δP

The reaction timescale:

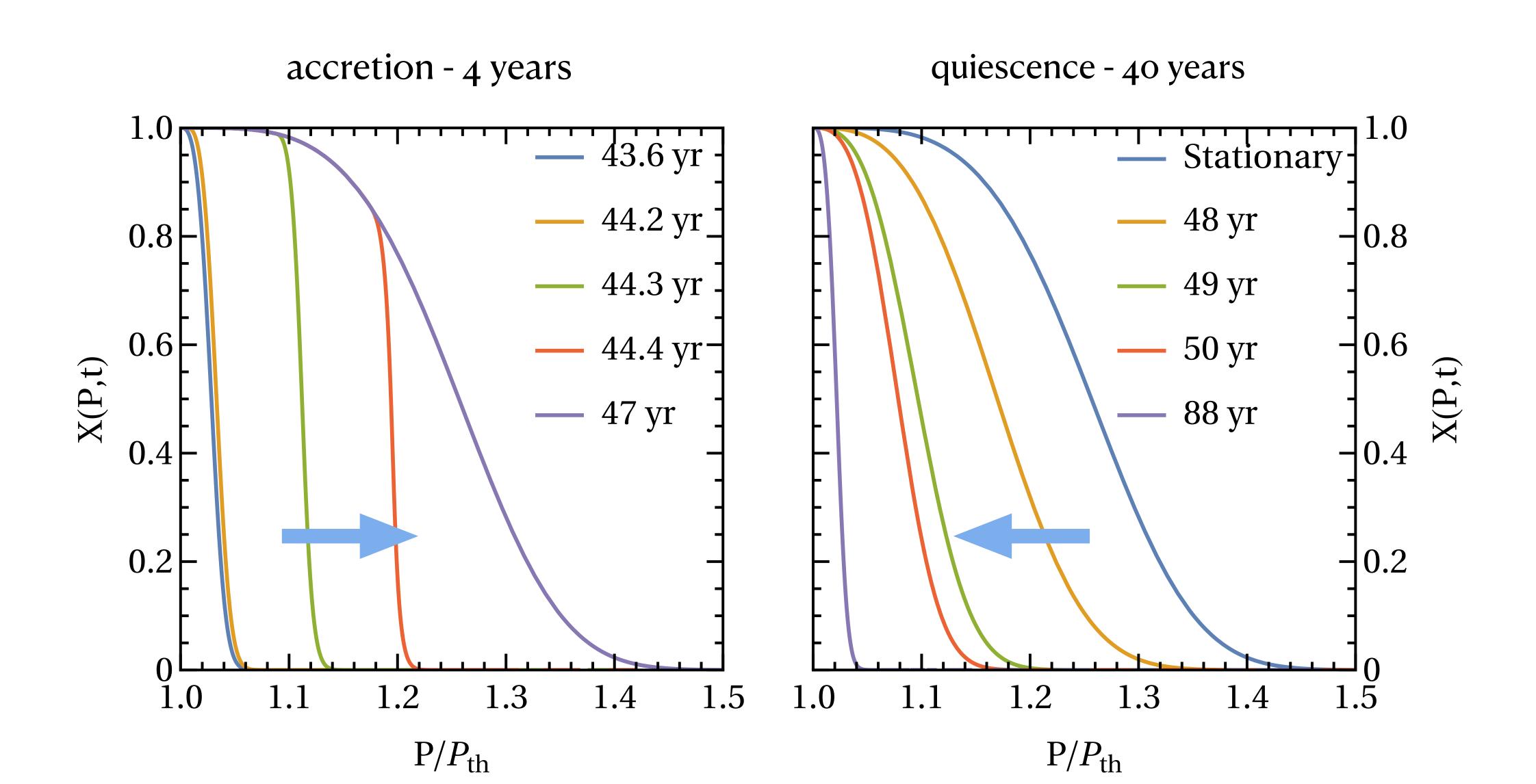
$$\tau_{react} = \tau_{react}(\delta P, X; (N, Z))$$

- decreasing function of δP



$$\partial_{\tau}(nX) + \partial_{z}(vnX) = -nX\mathscr{R}_{ec}.$$

Evolution of the reaction layer



Conclusions

- The heat sources are time- and pressure- dependent
- The total heat release is a function of astrophysical (active and quiescent time) and microscopic (reaction rate) parameters.
- These parameters should be considered separately and individually for each reaction layer.
- Thickness of the reaction layers: $\delta P = 10^{-7} \text{ MeV/fm}^3 10 \text{ cm} 10 \text{ m}$
- Shell thickness (between reaction layers) $\Delta P \sim 10^{-5} 10^{-4} \; \mathrm{MeV/fm^3}$
- Energy release larger than in the case of instantaneous reaction (by $\sim 20\,\%$)